

COMMISSION OF THE EUROPEAN COMMUNITIES

COSTANZA (R,Z)

by

E. VINCENTI and A. CLUSAZ

1971



Joint Nuclear Research Centre Ispra Establishment - Italy

Nuclear Studies Division

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Fuel elements of prismatic shape are considered. Up to 6 groups of delayed neutrons are allowed.

Temperature feedback of fuel and coolant are considered independently and affect the nuclear constants. Control rod movement or diffused poison

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Inlet temperatures and mass flow in the coolant channels may be varied according to any externally given time table.

Before starting the transient initial calculations put the reactor in criticality conditions at a given power level in equilibrium with the temperature feedback.

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ABSTRACT

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KEYWORDS

FORTRAN 2-DIMENSIONAL CALCULATIONS CYLINDERS PROGRAMMING NEUTRON DIFFUSION EQUATION HEAT TRANSFER CANNING

COOLANTS FUEL PINS FORCED CONVECTION TEMPERATURE COEFFICIENT DOPPLER REACTIVITY POWER COEFFICIENT CRITICALITY

<u>Contents</u>:

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Purpose	page 5	
Subdivision of the reactor	" 5	
Control rods	" 9	
Neutronic initialization	" 10	
Criticality search	" 10	
Dynamic Calculations	" 11	
Method	" 13	
Thermal calculation	" 17	
Temperature reaction	" 19	
Programme composition	" 21	
First overlay	" 21	
Second overlay	" 22	
Input instructions	" 27	
Restart and interrupted calculation	" 44	
Sample problem input	" 48	
Sample problem output	" 57	
Appendix A	" 71	
Appendix B	" 76	

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Purpose *)

The code COSTANZA (R,Z) treats the dynamics of a cylindrical nuclear reactor. Two geometrical dymensions: radius R and hight Z are considered, in the supposition that there is an axial symmetry.

The nuclear time dependent diffusion equations, in two energy groups are solved numerically with the finitedifference approxiamtion.

The heat transfer equation from heating element to coolant in the channel are also solved numerically. The geometry of the heating element can be of any type, because the element is characterized by thermal resistances and thermal capacities given in input. Flow rates and inlet temperatures of coolant are externally imposed as input function of time. The coolant is a single phase fluid which may be liquid or gas, and is not supposed to boil or condense during the transient.

This code has been developed to treat accidents of intermediate gravity in which detailed spatial aspects of the phenomena are essential, but which are not so grave as to produce alterations or damages of the core lattice, or changements of the state in the coolant.

Subdivision of the reactor

The reactor is supposed to be cylindrical with axial symmetry. Therefore only a vertical section of it

^{*)} Manuscript received on March 15, 1971

will be considered. (Fig.1). This section is a rectangular surface limited by the axis of the cylinder at one side and by the extrapolated boundary of the reactor at the other three sides. This area is subdivided with a cartesian pointlattice. In the radial direction the points may be arbitrarily spaced. In the axial direction the points are located at equal intervals within the core and they may be arbitrarily distributed in the upper and lower reflector. Within the core the axial point lattice is the same for neutronic and thermal calculations.

The outer points are located at the extrapolate boundaries of the reactor. The first two vertical arrays of points are located simmetrically to the axis of the cylinder to obtain the condition of zero space derivative at the axis. The lattice may have a maximum of 400 points, product of the radial and axial points. The radial and axial points cannot be more than 40 each. In Fig. 1 is a point-lattice of 12 axial x 10 radial = 120 points. The core may be subdivided in a maximum number of 30 regions having different core compositions; their boundaries are defined by points of the lattice. The initial criticality condition may be obtained by diffused poison or by the insertion of control rods.

Each region contains a uniformly distributed poison, whose concentration may be varied for criticality search to obtain the initial equilibrium condition, but will remain constant during the transient. The core is also subdivided in vertical control zones. There can be a maximum number of 10 control zones. They have the same height of the reactor, including top and bottom reflectors if existing. Each zone contains control rods which move in vertical direction independently from the rods of the other zones. Each zone is also characterized by a certain type of cooling channel which may be different from those of the other zones. The coolant and the rods may enter into the core from apposite sides (rod from top and coolant from bottom or viceversa) or they may enter from the same side. Conventionally we assume that the coolant inlet is at the top and the points of the lattice are numerated starting from the top left along the radii down to the bottom right. The entrance side of the rods will be indicated in input. The vertical boundaries of the control zones are defined by points of the lattice. They may not coincide with the vertical boundaries of the composition regions.

Fig. 1 indicates 5 regions, the first of which includes the top, bottom and radial reflector considered all of the same composition in this case; regions 2,3,4,5 belong to the core. Two vertical central zones (I,II) contain the rods.

- 7 -



8 axial points in the active part of the core

- 8 -

Control rods

In each control zone the rods move as a bank all together. They are represented by a diffused equivalent poison. The fast and thermal macroscopic cross sections of this equivalent poison are given in input. The poisoned part of the zone is limited

by a movable boundary A-B of Fig. 2. When the movable boundary lay between two points of the vertical array, a linearly interpolated value of poison is assigned to the space between the two points.



The initial insertion of the rods in each zone is given in input. However, in all the zones choosen for criticality search, the depth of insertion will be determined by the programme and will be the same for all these zones.

In all other zones not choosen for criticality search; the rods maintain their initial insertion. During the transient the rods can be moved in any zone arbitrarily chosen. The movement will be specified by a tabulation of insertions versus times given.in input for each chosen zone separately. For intermediate times a linear interpolation will be made. To obtain a step-wise movement two tabulated values must be given having times equal to $n \cdot \Delta l$ and $(n+1) \cdot \Delta l$. If the time of the transient overtakes the last tabulated time, the rods will maintain the last tabulated position.

In all other zones the rods will maintain their initial positions or the positions determined by the criticality search.

Neutronic initialization

The flux calculation starts with flat fluxes. A fixed number of calculations are made (Sub. INIZ) and at every step the fluxes are normalized with the factor

$$F = \frac{\text{Total Power}}{\int_{cone} \mathcal{E}\left[\sum_{finit}^{finit} \Psi + \overline{\sum}_{finim}^{thin.} \Psi\right] dv}$$

in order to obtain a wanted value of total power. Although the reactor is not critical the distribution of fluxes, thus obtained, is very close to the static distribution.

Criticality search

The criticality search can be made either by varying the concentration of poison in any wanted composition region or by varying the depth of insertion of the rods in any wanted vertical control zone. The calculation is repeated until the reciprocal of the period of the average thermal flux differs from zero by less than a fixed amount. To accelerate the convergence, only prompt neutrons are considered during this process. The delayed neutron precursor concentrations will be calculated at their equilibrium values once obtained the critical condition.

After every iteration the fluxes are normalized at a wanted power and the temperatures and corresponding feed back on neutronic are calculated. Once obtained the equilibrium condition at the wanted power the physical constants of the regions may be different from their initial values given in input.

Dynamic Calculations

Starting from the initial equilibrium conditions the code calculates the dynamic behaviour of the reactor in time.

The calculations of the flux distributions and delayed neutrons precursor concentrations will be repeated at every time step of constant length Δt . The thermohydraulic calculations and the corresponding temperature reactions are calculated at every n time steps, n to be specified in input. During the transient the power level and flux distributions evolve freely according to the reactivity introduced by any perturbation and to the temperature feed-back. The perturbations may be given by varying the depth of insertion of the rods, or the mass flow, or the inlet temperature in any wanted zone independently. The variation of these parameters as functions of time are given in input in tabular form. A convenient value of the time interval Δt is given in Appendix B as ratio $\Delta t/T$, where T is the shortest reactor period in the transient to be studied.

The maximum increase of power and the admissible error are considered. It is easy to have an a priori estimate of these values. Method

The two-groups time-dependent diffusion equations and the equation of the delayed neutron precursors are:

- 1) $D_t \nabla^2 \psi \overline{\Sigma}_R \psi + (1 \beta) (\Sigma_{f_1} \psi + \overline{\Sigma}_{f_2} \psi) + \sum_{i=1}^N \lambda_i C_i = \frac{1}{w} \frac{\delta \psi}{\delta t}$
- 2) $D_2 \nabla^2 \varphi \overline{\Sigma}_a \varphi + \overline{\Sigma}_h \cdot p \cdot \psi = \frac{1}{\nabla} \frac{\delta \varphi}{\delta t}$
- 3) $\frac{SC_i}{\delta t} = \beta_i \left(v \sum_{q_i} \psi + v \overline{L}_{q_2} \psi \right) \lambda_i C_i$

As the reactor is cylindrical, with axial symmetry, a calculation of the fluxes and other physical magnitudes made in a vertical section delimited by the axis and outer boundaries, as shown in fig.3, describes completely the behaviour of the whole reactor.

This section is covered with a point lattice. The reactor volume is subdivided into cylindrical annuli each of which is delimited by four points of the lattice.

- j = is the radial index increasing from the axis outwards
- i = is the axial index increasing from the upper boundary downwards

fig.4

(*-1,1)

The physical constants of the equations 1) 2) 3), generally indicated with Aij, may have different values for each annular part of the space, but are constant within the annulus. There is a correspondabce of indexes between the part of the space and the lower left of the four points which delimit it.

The equations 1) 2) 3) are integrated over the volume corresponding to the point (i,j) which is subdivided into four sub-volumes

see



Each sub-volume belongs to a different part of space, which may have different physical constants. Thus for every point of the lattice, the integrated equations are averaged over the surrounding parts of space. This is evidently significant for points on the interferences of different regions, but it is also significant for the internal points of a region, because the variations of the physical constants due to the temperature reaction are calculated pointwise. By this method the formal expression of the integrated equations results to be the same everywhere with a considerable simplification of the programme.

For the details of the integration see Appendix A.

The general form of the integrated equation for the generic point (i,j) is:

This is the eq. corresponding to the thermal flux and contains, in the known term of the right side the source term $b_{ij} \ \psi_{ij}$ with the fast flux ψ_{ij} . The equation corresponding to the fast flux ψ is formally the same and contains a source term with the thermal flux φ .

There are two systems of n linear equations, one for each energy group; n is the number of the points of the lattice.

Note

The boundary condition $\frac{d\varphi}{d\tau} = 0$ at the axis of the cylindrical core, is obtained by putting all the points (i,1) and (i,2) of the first two vertical columns, symmetrically to the axis. i.e. r(1) = -r(2) (r(j) being the radial distance of the point (i,j) from the axis), and by imposing $\varphi_{d,1} = \varphi_{d,2}$. This is obtained by putting $\beta_{d,1} = 0$. In fact the equation corresponding to the first point of each radius is of the type:

8)
$$-\alpha_{i,2} \varphi_{i,2} - \beta_{i,2} \varphi_{i,4} - \delta_{i,2} \varphi_{i+1,2} - \delta_{i,2} \varphi_{i,3} + \epsilon_{i,2} \varphi_{i,2} = b_{i,2} \varphi_{i,2} + c_{i,2}$$

if $\varphi_{i,2} = \varphi_{i,2}$ then the equation can be written:

9) $-d_{i2} \varphi_{i2} - g_{i2} \varphi_{i+1,2} - \delta_{i2} \varphi_{i,3} + \left[g_{i,2} - \beta_{i2} \right] \varphi_{i,2} = b_{i2} \varphi_{i2} + c_{i2}$ as β_{i2} is also contained in g_{i2} with the + sign the equation 9) can be obtained from 8) by simply putting $\beta_{i,2} = 0$. The outer boundary conditions are obtained by imposing $\varphi = 0$ and $\psi = 0$ in all the points of the outer boundary. In a first version of the code we have used the Gauss-Seidel iterative method with overrelaxation factor for the solution of each groups separately, using than the obtained solution as source for the other groups and iterating between the groups.

This way of proceeding has been abandoned for a much faster method in part direct and in part iterative. In eq. 4) the fluxes $\psi_{i-1,j}$ and $\psi_{i+1,j}$ correspond to points belonging to the radius (i-1) above and to the radius (i+1) below the point (i,j).

Let us consider the values $\varphi_{i+1,j}$ and $\varphi_{i+1,j}$ as already known. All the equations of the i radius form a system of JMAX equations (JMAX = number of points of the radius), which has a tridiagonal matrix of coefficients. In fact each equation contains only three incognita. This could be solved with the well known method of forward elimination and backward substitution, for each group separately. However by intercalating the equation of the fast with those of the thermal group and rearranging the terms, it is possible to obtain a system of 2 x (IMAX) equations for both groups, with a pentadiagonal matrix of coefficients. This can be partitioned into (2x2) submatrices and considered as a tridiagonal matrix with entries consisting of (2x2) submatrices. The method of forward elimination on backward substitution can be applied operating on these (2x2) submatrices and the direct solution can be obtained for both groups at the same time. (See: Finite Difference Method for solving the Spatio-Temporal Diffusion Equation in the Two-group Approximation - EUR 596 by B.Monterosso and E.Vincenti).

- 16 -

All the point lattice has to be scanned by repeating this direct calculation at each radius. For $\varphi_{i+i,j}$ will be used the value obtained in the calculation of the $(i-1)^{th}$ radius. In the actual iteration, and for $\varphi_{i+i,j}$ will be used the value obtained in the preceeding iteration. The scanning of the lattice will be repeated alternatively downard and upward until convergence is reached. This method which is direct in radial direction and iterative in axial direction eliminates the iterations between the groups and in the normal cases reduces the computing time of a factor 3 compared to the usual overrelaxation method.

Thermal calculation

The reactor is subdivided in a number of vertical zones (max. 10). Each zone contains many vertical cooling channels of the same type. Channels belonging to different zones may however be of different types. For channel is meant here a fuel rod with the corresponding coolant. The channels go through the reactor from top to bottom including the reflectors, but only the part contained in the active core is considered for thermal calculation. The part is subdivided with a point-lattice of constant mesh $\Delta \geq$, which coincide; with the axial lattice for neutronic calculation.

To determine the power developped in a fuel rod, in any zone, in a part Δ Z between the (i-1)th and the (i)th level, the power is integrated in the volume A (Fig.6) of the zone between (i-1) and (i) and this is divided by the number of rods of the zone.



number of rods in the zone

The heat is generated in the rod and propagates to the coolant into which it is injected at the point (i). The fuel rod is subdivided in a number of concentrical shells (5 10). The relative distribution of power among the shells is given in input and remains constant during the transient, although the total power varies with time. The time-dependent heat balance equation, transformed with the finite-difference method, is solved in implicit way at each level (i) and gives the temperature distribution at a given instant in the shells of the fuel and in the coolant. The results obtained at level (i) will be used as input for the calculation at level (i+1). A complete description of this procedure is in the report EUR 3171 by A.Agazzi, G.Forti, E.Vincenti, pages 8 14.

Temperature reaction

Eight physical constants are allowed to vary with the temperature, i.e.:

this temperature dependence has been expressed in the form of a Taylar expansion:

 $C = C_o + \alpha \left(T_f - T_{f_o} \right) + \beta \left(T_{au} - T_{au_o} \right) + \gamma \left(T_f - T_{f_o} \right)^2 + \beta \left(T_{au} - T_{au_o} \right)^2$

where

 C_o = initial value of the physical constant

 $\alpha_{,\beta}, \gamma_{,\delta} = \text{temperature coefficients}$ $T_{f_{\nu}}$ and $T_{m_{\nu}} = \text{Reference temperatures}$ of fuel and moderator corresponding to the initial value C_{ν} ⁽¹⁾ All these magnitutes are given in input region wise.

I MARCH THE MART

 T_{f} and T_{m} = New values of the fuel and moderator temperatures. (1)

(1) <u>Note</u> This version of the code is made for pressurized water reactors where the moderator and coolant are at the same temperature. Another version of the code has been made for HTG pebble-bed reactors, where the moderator (graphite) is contained in the fuel element mixed with the fuel. Although all other values, including the reference temperatures T_{fo} and T_{mo} , are given for each composition region, the temperatures T_f and T_m are calculated in each vertical control zone. The calculation is made at every thermal time step. T_f and T_m are values of temperatures averaged over that part of the radius within the control zone, they have however distinct values at every axial level. The temperature reaction is therefore made pointwise in axial direction, and zonewise in radial direction.

Programme Composition

The code consists of a Main and 19 Subroutines of which 5 belong to the first Overlay and the remaining 14 belong to the second Overlay.

The 5 routines of the first group are called only once. They calculate all the coefficients which will not be changed anymore and will be transferred to the second overlay in order to make available the rest of the memory.

First overlay :

- 1) MAIN Calls the first routines of the first group and PRINC 2, control routine of the second group.
- 2) INPUT Reads the vector DATA from cards or tape.
- 3) PRINC It takes the data which are read region-wise in Input and transform^S them into point-wise data. Calls GEOM and DCAN.
- 4) DCAN Calculates all the thermohydraulic coefficients which will remain unchanged during the transient.
- 5) GEOM Calculates all the coefficient depending from the geometry of the core and the point lattice.
- 6) AZER 1 Called from Input, before reading the data. Puts equal zero all the commons of the first Overlay.

7) PRINC-2 - In its first part, before label 1000, this routine determines the initial conditions. It calls AZER 2 to put equal zero all the commons belonging only to the second Overlay. Puts the control rod in the initial position. Calls INIZ to give a first approximate distribution of the fluxes. Calls CRITIC to calculate the initial equilibrium distributions of fluxes and temp@ratures.

In the <u>second part</u> after the label 1000 it determines the transient behaviour.

At every time step it calls BARRE to determine the rod position; it calls Flussi to calculate the flux distribution; optionally it calculates the accumulated energy per cm³ at every point; it calculates the delayed neutrons precursors concentrations at every point.

At every KPC (see input key) time steps it calls the thermal routines CANPALL and CØVAR to determine the temperature reactions.

Calls STAMPA to print according to the orders given in input.

Controls the restart operations to store on tape the results at the end of a calculation and to start from tape an interrupted calculation (see Restart Instructions).

8) AZER 2 - Puts equal zero all the commons belonging only to the second Overlay.

- 9) INIZ It calls FLUSSI a number of times specified in input. Starting from and initial flat distribution determines a first approximate distribution of the fluxes. In these calculations the delayed neutrons are not considered. It normalizes the fluxes at a given initial power level. Uses a neutronic time step DELTI which may be different from DELTC used in CRITIC, and from DELT used for the transient calculation.
- 10) FLUSSI It solves directly along each radius the two-group system of linear equations. (see Method). It scans axially the point lattice, in alternate directions, until the residuals of each group

$$\frac{\sum_{i=1}^{N} |\Delta \psi_i|}{\sum_{i=1}^{N} |\Psi_i|} < EP1 \quad and \quad \frac{\sum_{i=1}^{N} |\Delta \psi_i|}{\sum_{i=1}^{N} |\Psi_i|} < EP2$$

are smaller of the quantities EP1 and EP2 given in input.

It calculates the fast and thermal flux average values and the reciprocal of the period.

11) MAT1 - It calculates the coefficients of the system of linear eq. to be solved by Flussi. This calculation is made using the invariant coefficients transfered from the first overlay and variable physical magnitudes depending from the temperature reaction and the rod position. This routine will be called at every thermal calculation, from Princ 2 and from Critic. When only the poison concentration is varied it is enough to call MAT2, because, as the routine contains the Entry Mat 2, the first part of it will be bypassed.

- 12) STAMPA It prints the output. Any wanted result may be printed in several optional types of print (see input instructions).
- 13) CRITC searches the criticality by varying iteratively either the concentration of poison in any wanted region, or the insertion of the rod in any wanted zone. In the first part of this iterative method from two values of poison and the corresponding values of reciprocal period, it determines, by extrapolation, the next value of poison and so on. When the reciprocal period is below a given value it variates the poison at each iteration of one thausand of the last concentration, until the convergency criterium is satisfied or the reciprocal of period changes of sign. The same proceeding is followed when the criticality is searched with rod movement. After a first part with the extrapolation method, the rod are displaced of 1 millimeter at each iteration. In both cases at each iteration the thermohydraulic routines are called to determine the temperature reaction. The calculation is made without considering the delayed neutrons.
- 14) BARRE This routine calculates the new coefficients of the system of linear equations only where a new condition of poison was determined by the displacement of the rods e.g. see M and N of Fig. 7.... The recursive coefficients of the forward elimination

and backword substitution method however are calculated by BARRE in all the points of the region A B C D of Fig. 7 determined by the rods displacement and the outer boundary.



It is called from Critic and Princ 2. When called from Princ 2 it determines the rod position by interpolation of the tabulation given in input.

- 15) TEST This is a dummy routine which contains all the commons and may be used to introduce any new statement in the programme. It is called at every time step from Princ 2.
- 16) CANPAL This routine calculates from each channel the temperature distribution in the fuel rod and in the coolant at all axial levels (see thermal calculation). It is called during the transient from Princ 2 at every KPC time step, and from Critic to determine the equilibrium initial condition.

- 17) VINIZ Determines for each channel the coolant inlet temperatures and mass flow at every time step, interpolating the values given in input as function of time in tabulated form.
- 18) INTEGR Calculates for each channel independently at every axial level the fuel and coolant average temperatures.
- 19) COVAR calculates the new values of the physical constants given by Integr. (See temperature reaction).
- 20) RESTAR stores on tape the results at a given time. It restarts an interrumpted calculation, reading the results stored on tape (See restant instructions).

Input Instructions

A title card is the first input card of each problem. A vector of 10000 memory positions : DATA (1)..... DATA (10000) contains all the data.

Since entire groups of memory positions are zero, it is possible to read separate sets of significant data; each set of data must be preceeded by a card containing the integers Ki_1 , Ki_2 , Ki_3 , written in Format (216, 112) where Ki_1 is normally zero, Ki_2 and Ki_3 are the indexes of the first and last datum of the set.

Only in the last set of a problem Ki₁ must be any positive integer number. The data are in floating form, six per card (Format 6 E 12.8).

Any number of problems may be run in sequence; for all the subsequent problem only the data different from the corresponding data of the first problem need to be given. The vector DATA of the first problem must be memorized on an auxiliary tape. (See Restart Instructions).

Input Key

<u>Title card</u>: Columns 1-2-3 if left blank the data are read from cards and the results will not be stored on tape. Numbers in these columns correspond to store and restart (see Restart Instructions). Columns 4-5-6. If only one problem is treated they must contain any positive integer. If a series of problems are to be treated in the same run, only the last title card of the last problem must have any positive integer, in the preceeding title cards these columns must be blank. Columns 7-70 may contain any alphanumerical information which will be printed in output.

	Nu	DATA umber	Variable Name	DESCRIPTION	Allowable Values	NOTES
		(1	TEMAC	Computer time (min.)		Only if the first three columns of the title card contain a number see Restart Instructions.
•		2	DELT	Neutronic time step Δ t for transient calculation (sec.)		· · ·
	card	3	DELTC	Neutronic time step ムt for criticality search (sec.)		
	st	4	IMAX	Number of axial mesh points	\$ 40	(IMAX * .TMAX) < 400
	Fir	5 .	JMAX	" "radial " "	≼ 40	
		6	NREG	Number of regions with different composition	\$30	
				· · · · · · · · · · · · · · · · · · ·		
	1	(7	NRIT	Number of delayed neutron groups	< 6	
		8	ITMAX	Maximum allowable number of axial iterations	suggested value = 30	*
	card	9	IDST	Number of initial iterations before criticality search	Suggested value = 20	
1	puq	10		blank		
	Sec	11		blank		
		12	DELTI	Neutronic time step Δt for the initial iterations before criticality		

l	DATA Number	Variable Name	DESCRIPTION	Allowable Values	NOTES
	(13		Blank		
	14	EP1	Convergence criterium for the fast group	s.v.=0,01	
יס	15	EP2	Same for the thermal group	s.v.=0,01	
car	16	SI	Initial reactor power (Watts)		
hird	17	KPC	The thermal calculation is done every KPC neutronic steps		
	18	NCAN	Number of radial zones containing channels		
			•		ı 29
	19		Blank		1
	20		Blank		
	21	KTE	If KTE > 0 Subroutine TEST will be called.		
	22	ICA	Axial index of the points laying on the upper core boundary. (Reflector excluded)		
	23	ICB	Same for the lower boundary		
	24	JCI	Radial index of the points laying on the inner core boundary		Usually the central part of the reactor belongs
	25	JCE	Radial index of the points laying on the outer boundary of the core. (Radial reflector excluded)		to the core, then $\underline{JCI=2}$

DATA Number	V ari able Name	DESCRIPTION	Allowable values	NOTES	
26	W	Fast group neutron velocity (sec)			-
27	v	Thermal " " "			
28	SPRG	First guess of control poison concentration		Blank if cr itically search is made with	
29	DAPF	Convergence criterion for reciprocal of period in criticality search	s.v.=0.1	rod movement	
30	LF	Maximum allowable number of criticality iterations	s.v.=30		
31	IDIR	<pre>{ = 1 if the control rods enter</pre>			- 30
		= 2 if they enter at the opposite side		i	1
32	IBCR	<pre>{ = 0 criticality search made with diffused poison 1 criticality search made with</pre>	• •		
		rod movement			
33	NZB	Number of the vertical zones	≤1 0	The vertical zones have the same height of the reactor. Each zone is	

30 L

characterized by the same rod depth of insertion and by the same typical channel.

DATA "Number"	Variable ^{Ab} Name	DESCRIPTION	Allowable values	NOTES	
34	SPB1	Microscopic control poison (barn) cross section for the fast groups.		they are multiplied by SPR in each region to obtain the correspondin macroscopic cross section	
35	SPB2	Same for the thermal group			obtain the correspond
201-206	BETA	eta_{ι} delayed neutronsyield per fission			
211-216	DL	λ_i delayed neutrons precursors decay constants			
231	INTG	<pre>{ = 1 The power will be integrated in the time at every mesh point = 0 No integration will be effectuated</pre>			
232	NU SO	Number of thresholds of integrated energies	\$ 10	only if INTG = 1	
233-242	SØLEN	Thresholds of the integrated energies in increasing order from the lower to the higher	[Joule] cm ³	(See printing order N° 8 and <u>9</u>)	
. 301	I1	Index of the upper boundary of the first region		This group of data for energy regions is	
302	12	Index of the lower boundary of the first region		included in one card. For the region bounder by the axis of culind	
303	J1	Index of the left (inner) boundary of the first region	J1=1	J1=1	
304	J2	Index of the right (outer) boundary of the first region			
305	-	blank			
306	-	blank			
307-312		for the second region			
375 - 480		for the 30th region			

DATA Number	Variable Name	DESCRIPTION	Allowable values	NOTES	
	· · · · · · · · · ·		<u> </u>		
1001	D1	Diffusion coefficient fast group	· · · · · · · · · · · · · · · · · · ·		
1002	SR	Removal cross section			
1003	SF1		or		
1004	P	Resonance escape probability	d f		
1005	-	blank	car re		
1006	_	blank	ach	.	
1007 - 1012 1175 - 1180	• • • • • •	Second region 30th region	C 0		
	•				
		•			
1601	D2	Diffusion coefficient thermal group			
1602	SA	thermal absorption cross-section	1		
1603	SF2		<u>ب</u>		
1604	SPR	Poison concentration; it is multi- plied by SPB1=Data (34) and SPB2=Data (35)	card fo region	If IBCR = 0 SPR will	
1605 1606 1607-1612	- -	blank blank Second region	One c each	be varied to obtain criticality in the control regions	
 DATA Number	Variable Name	DESCRIPTION	Allowable values	NOTES	
---------------------------	------------------	---	----------------------	---	--------------
2201	EQUi1	Thermal equivalent of fission (fast)			
2202	EQU12	Thermal equivalent of fission (thermal) [Watt-sec]			
2203	FNU	$oldsymbol{\gamma}$ Neutronyield per fission			
2204	KV	<pre>KV = order number of the region, then</pre>		Only if IBCR=0	
2205	TRU	The initial values of the physical			
2206	TMU	constants given in input were calculated according to a fuel temperature = TRU and a moderator temperature = TMU for each composition region.			1 33 1
2207-2212		for the second region			
2375-2380		for the 30 th region			
2801	KB	<pre>KB= order number of the zone, this zone contains control rods</pre>			
	•	KB=0, no rods		See note	
2802	KBC	<pre>KBC= order number of the zone, the rods of the zone are used for criticality search kEC= zero, no movement of rod for criticality search</pre>	Only if IBCR=1	of Data (33) - The order number of zone, must increase from left to right	the

DATA Number	Variable Name	DESCRIPTION	Allowable values	NOTES	
2803	J1B	Index of the left (inner) boundary of the zone		For the region bounded by the axis of cylinder	r
2804	J2B	Index of the right (outer) boundary of the zone		J1B = 1	
2805	KCA	KCA= order number of zone, the zone contains channels			
-		kcA = zero, no channel			
2806	_	blank			
2807-2812		Second zone			
2855-2860		10 th zone		1	
		•		34 44 1	
2861	VEBA1	Macroscopic cross section of the rod equivalent poison (fast)			
2 86 2 ·	VEBA2	Macroscopic cross section of the rod equivalent poison (thermal)			
2863	-	blank			
2864	_	blank			
2865	_	blank			
2866	-	bl a nk			
2 867 - 2872		Second zone			
2915-2 920		10 th zone		· · ·	

DATA Number	Variable Name	DESCRIPTION	Allowable values .	NOTES
3001-3030	TOBA	Times of the rod movement tabulation for the first zone sec	DATA (3001) mus	t always be zero
3031-3060	PRIN	Corresponding depth of rod insertion in the tabulation for the first zone cm		
3061-3090		Same for the second zone	DATA (3061) = 0	
3091-3120		etc.	DATA (3541) = 0	۱ د
4001 - 4040	X(J)	Distance of mesh points from the axis		
4101-4140	Y(I)	Distance of mesh points from the upper boundary.		The upper boundary is for definition the side from which the coplant enters.

.

The first distance is Y(1) = 0.0

DATA Number	Variable Name	DESCRIPTION	Allowable values	NOTES
		<u>CHANNEL DATA</u>		
	•			
5001	DZ	Constant axial mesh distance cm		,
5002	SEZ.	Section of channel cm ² (Coolant corresponding to a fuel rod)		
5003	SC	Coolant specific heat <u>Joule</u> g °C		
5004	RO	Coolant density g/cm ³		
5005		blank		
		hlank		

DATA Number	Variable Name	DESCRIPTION	Allowable values	NOTES
5007	N: (KC)	Number of shells in the fuel rod	N1 9	• • •
5008	PALN (KC)	Number of fuel rods in the zone KC		
5009		blank		
5010	ØSE1	Outer surface of a part of rod of height DZ		ч.
5011		blank		•
50 12	-	blank		1
		·		3 7 1
5013-5018		These six DATA contained in one card correspond to the first shell of the rod.		
5013	RES1	Thermal resistence between the shell and the next one [°C/Watt]		
5014	CPT1	Thermal capacity of the shell [Joules/°C]		· · · · ·
5015	VOS1	Shell volume [cm ³]		
5016	PZO1	Power fraction corresponding to the shell		Relative values are significant. Normali- zation is performed by the code.

•

DATA Number	Variable Name	DESCRIPTION	Allowable values	NOTES
501 <u>9-</u> 5024		Same data for the second shell etc		
5061-5066		Same data for the nineth shell		
		• <u></u>		
5190	VITE	Inlet temperature of coolant at equilibrium initial condi= tions (Put zero if inlet tempe= rature is tabulated as function of time) [°C]		י 38 י
5191	STEPT	Step of coolant inlet temperature		Only if VITE $\neq 0$
5192	RAMPT	Value of dT/a tfor ramps in coolant inlet temperature [°C/sec]		tt .
5193	VIFO	Iniatial value of coclant mass- flow (put zero if tabulated) [g/sec]		
5194	STEF PO	Step of coolant mass flow		Only if VIPO \neq 0
5195	RAME C	Value of $\frac{dW}{dr}$ for ramp in mass flow [g/sec ²]		1 9
		•		

5201-5329 SECOND CHANNEL DATA 6801-6999 TENTH CHANNEL DATA NOTE: All the cards must be repeated for every channel. The order number of the e channel must be the same of the order number of the radial zone in which th is contained.	existent 1e channel	
6801-6999 TENTH CHANNEL DATA NOTE: All the cards must be repeated for every channel. The order number of the echannel must be the same of the order number of the radial zone in which the is contained.	existent 1e channel	
NOTE: All the cards must be repeated for every channel. The order number of the e channel must be the same of the order number of the radial zone in which the is contained.	existent 1e channel	
(Ex. Suppose that zone one and three contain channel and zone two do not co channel, then the channel of zone three must be considered as the third and are in the group Data (5401) (5599).)	ontain any l _i its data	
7001-7060 Tabulation of inlet temperatures and mass flow for the <u>first</u> <u>channel</u> . Each point of the tabu- lation for both magnitudes is contained in one card (4 data), as follows:		
7001 TOTE Time of the temperature tabulation (Zero in the first card)		
7002 TENT Value of the corresponding tempera- ture		
7003 TOPO Time of the mass-flow tabulation (zero in the first card)		
7004 PENT Value of the corresponding mass flow	· :	
7005 PENT blank		
7006 PENT blank		

DATA Number	Variable Name	DESCRIPTION	Allowable values	NOTES
7007-7012		Second card of the tabultation etc.		
7061-7120		Same for the second channel etc.		
7541-7600		Same for the 10 th channel		

TEMPERATURE COEFFICIENTS ($8001 \rightarrow 8960$)

The temperature dependence of the physical constants is expressed in the form of a Taylor expansion,

$$C = C_{c} + \alpha_{1} \left(T_{f} - T_{f} \right) + \alpha_{2} \left(T_{m} - T_{m} \right) + \alpha_{3} \left(T_{f} - T_{f} \right)^{2} + \alpha_{4} \left(T_{m} - T_{m} \right)^{2}$$

The variable physical constants are 8 in the following order :

1
$$D1 = D_{4}$$

2 $SR = \sum_{R}$
3 $SF1 = \nu \sum_{f_{4}}$
4 $P = f_{2}$
5 $D2 = D_{2}$
6 $SA = \sum_{a_{2}}$
7 $SF2 = \nu \sum_{f_{2}}$
8 $SPR = \sum_{a_{2}}$ poison

40 -

The algoritm giving the position of the temperature coefficient in the reactor DATA is as follows:

$$CTN(K, L, N) = DATA[8000 + 32*(N-1) + 8*(L-1) + K]$$

DATA Numb er	Variable Name	DESCRIPTION	values	NOTES
			······································	

where K(1,...8) = order number of the physical constant

L $(1, \ldots 4)$ = order number of α_i in the Taylor series

N $(1, \ldots 30)$ = order number of the composition region

41 -

PRINTING ORDERS

9901–9960	It is possible to give in input up to 5 different printing orders. Each order is contained in two cards (12 data)
9901-9912	First order
991 3 9924	Second "
	11 .
9949 - 9960	Fifth "
	In each group of 12 data (two cards) corresponding to an order the first datum = n1 indicates that the printing order is valid up to the $n1$ th time-step.

DATA Number	Variable Name	DESCRIPTION	Allowable values	NOTES
		The other data of the group $n_i \leq n_1$ (i= 2,3,12) indicate that the corresponding type of printing will be made at every n_i time-step. If some of the data = n_i are left blank the corresponding type of printing will not be made.	· · · · · · · · · · · · · · · · · · ·	
		`•		
	Ex. For the fi	rst order :		1 42 1
9901	n 1	The first printing order will be performed up to the n_1 th time step.	γ	
9902	n 2	Flux core average, reciprocal period and total power	ro	
9903	nz	Flux map	car	
9904	n 4	Map of rod insertions	n t	
9905	n 5		、 氏	
9906	n 6	Map of the concentrations of the delayed neutron precursors		
		•	· · · ·	

DATA Number	Variable Name	DESCRIPTION	Allowable values	NOTES
9907	n 7	Maximal fuel temperature, max. cladding temperature, max. heat flux, inlet and outlet temperature, mass flow.		
9908	n 8	Map of temperatures	гq	
9909	n 9	Percent of core (in volume) below a given integrated energy threshold	cond ca	•
9910	n1 0	Map of integrated energies	ບ ບ	ا 44
9911 9912		nothing		

It is possible to give as many order as wanted up to five. The calculation stops at the time-step indicated by the first datum =n1 of the last printing order.

Restart an interrupted calculation

Two dimensional calculations require long computer time. It may be interesting to have the possibility of restarting an interrupted calculation, without repeating it from the beginning. In particular it may be interesting to have the possibility to initiate again a calculation from the equilibrium condition, without repeating the criticality search. For this purpose it is necessary to store on tape the results already obtained.

It is possible after a restart to change the printing orders. The first card of the new printing orders must contain in the first twelve columns the number IT of the time steps at which the calculation was interrupted; this card is blank in all other columns. A second blank card must follow, then groups of two cards must follow containing the new printing orders. It is also possible to change the tabulations given in input for the rod movements, inlet temperatures, and mass flow. Of course only changements after the instant **i** of restart have a meaning.

The new tabulations must contain a first point at time zero, with any value of the variable, the second point will be at the instant t and the value of the variable (insertion of rod, inlet temperature, mass flow) must be that indicated in the output of the results which were stored on tape. From t on any new point may be introduced in the tabulation.

At any subsequent restart after having changed the original tabulation, it is necessary to use the cards of the new tabulation, otherwise the programme will calculate according

- 44 -

to the original one which remains stored on tape.

It is also possible in a restarted calculation to introduce a new value of the time step Data (2). In this case a new printing order must be given, as described in the preceeding page. The original vector Data is stored on tape, only the new changed values must be given by cards.

The orders of restart are given in the fisrt three columns of the title card:

Titel Card:

Blank	Ξ	results on tape	from card, i	no storing of
001	-	Input data read stored on tape	from cards,	three files are
		file 1	input data	
•		file 2	results of	criticality search
		file 3	results of	instant t.
		· .		

- 002 = Only if more than one problem is treated. Input data corresponding to the first problem are read from tape, file 1. Only those input data of the new problem which differ from those of the first problem are read from cards. At every new restart of the problem the cards corresponding to changement of the input data must be introduced, because on tape will always remain stored the input of the first problem.
- 003 = Restart from equilibrium conditions after criticality search, and store results at the end of calculation at time t.

004 005 = restart from results stored on file 006 (n-1) and store results on file n. 00n

-03-04-05-0n

Restart from the results stored in any file without destroying the subsequent files. This in case of temptative changements of rods, inlet temperature, mass flow tabulation and time step in any part of the transient.

First data card:

In addition to the titel card the first data card, containing DATA (1) to (6) must always be present. In DATA (1) = TEMAC is indicated the computer time (minutes). At every thermal calculation the programme will compare the remaining computer time with the computer time necessary to make a thermal time step. If it will not be possible to make the next thermal calculation, then the results will be stored on tape.

If no thermal calculations are made then the programme makes this check at every 100 steps.

Every DD statement corresponding to a file stored on tape must contain the parameter DISP = (NEW, DELETE), and every DD st. corresponding to a file read from tape must contain the parameter DISP = (\emptyset LD, DELETE). EURATOM - CETIS

DATA (Zones de 10 colonnes)

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LABEL = /(2), SLI DSNAME = EU27	472		C
DISP= NEW DELETE!			
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DISP=(NEW DELETE)			
IGS SYSIN DD *			

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- 49 -

VI - CET 002

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VI.CET 002

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- 53 -

VI-CET 002

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VI.CET 002



VI - CET 002

COSTANZA R.Z.

SAMPLE PROBLEM

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13 19 96		19		15 21	0.1000008-02	16 22	C.3CCOCCE 09	17	C.1C00005 03 0.1CC0CCE 02	18	C.200000E C1 C.2C0000E 01		
31	0.10000000	32)•)	33	0.2000009 01	28 34	0.2000005-23	35	0.8000005-23	30	0.3000000 02		
201	0.2190008-03	202	0.1116005-02	203	C.977CCC5-03	204	0.1687008-02	205	0.4550005-03	2 C 6	C.1950CCE-03		
211	0.1250008-01	212	0.3190006-01	213	0.12300000 00	214	0.312000£ CO	215	C.1150005-01	216	0.2740CCE 01		
231	0.1(0000E 01	232	3.1000000000000000000000000000000000000	233	0.1000008 02								
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1001 1007	C.137CCOD C1	1002	0.3100000 - 02	1003	0.0 0.0	1004	G.100000E 01	1005	C • C	1006	C • C	Input	57
	C.135000E 01 0.135000E 01	1014	1.335CC03-C2 0.335CC02-C2	1015 1021	C.C.C	1016	C.1CCCCCE 01 C.1CCCCCF C1	1017 1023		1018 1024	C.C Q.C	Data	Т
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1607	2. 977 2005 0C		0.4520008-02 0.4520008-02	1609	0.5200006-02	1510				1512			
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2201	C.C. C.317200E-10	2202	0.0	2203	0.C 0.2460C08_01	2204	C.1CCCCOB 01	2205	C.C. C.8C000CE 03	2206	0.C		
2213	0.317200E-10 C.3172C0E-1C	2214	0.3172005-10 0.3172005-10	2215 2221	0.246000E 01 0.2460C0E 01	2222	0.3C00C0E C1 C.4CCCCCE C1	2217	C.8C0000E 03 C.8CCCCCE 03	2218 2224	0.290000E 03 C.290000E 03		
2225	0.317200L-10	2220	J •317200L-1C	2227	0.24600DE 01	2228	0.5000000 01	2229	0.800005 03	2230	0.290000F 03		
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3001	0.C	3002	0.200003 CC										
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3091	: 〕	3092	0.2450000 (3	3093	C-6120CDE 03			<u>۲</u>
4001 4007	-0.3000001 02 0.2000000 03	4002 4003	0.300005 C2 3.2300001 C3	4003 4009	C.6CCOCOE 02 4004 0.240000E 03 4010	C.9CC0008 C2 4005 C.250000E D3	C.12000CE 03 4006	C.16CCOOE 03
$\begin{array}{c} 4101 \\ 4107 \end{array}$	0.0 0.344/000E 03	4102 4108	0.4200002 02 0.4200002 03	4103 4103	C.4CCCCOF 02 41C4 J.4960000 03 4110	0.11600000 03 4105 0.5720005 03 4111	0.1920005 03 4106 0.5920005 03 4112	C.2660CCF C3 0.612000E 03
5001 5007 5013 5025 5025 5031 5032 5031	C.76CC00E 02 C.5C0000E 01 0.775C00E-01 0.2335CCE-01 C.1347C00E-01 C.1393CCE-01 C.1393CCE-01 C.1393CCE-02 C.377C0CE-02	20000000000000000000000000000000000000	0.13CCC3C C1 0.2773CC3C C4 C.177CC0C-C1 0.377CC0C-C1 0.377CC0C-C1 0.377CC0C-C1 0.377CC0C-C1 0.377CC0C-C1 0.377CC0C-C1 0.377CC0C-C1 0.377CC0C-C1	5003 5005 5021 5021 5023 5033 5033 5033 5033	0.540CCO2 C1 5004 0.C 5010 0.11383E 02 5016 0.11383E 02 5022 0.11383E 02 5028 0.11383E 02 5028 0.119383E 02 5034 0.119383E 02 5040 0.125321E 02 5040	C.8C00CCE CC 50C5 0.286000H C3 5011 C.2000CCE CC 5017 C.2C00C0E CO 5023 C.2C00C0E CO 5023 C.2C00C0E CO 5025 C.2C00C0E CC 5035 C.2C00C0E CC 5035 C.2C00C0E CC 5035 C.2C00C0E CC 5035 C.2C00C0E CC 5035	0.0 50C6 0.0 5012 0.0 5012 0.C 5024 C.0 5036 C.0 5036 C.0 5036 C.0 5036 C.0 5048	
5130	0.2700008-03	5131	0.0	5132	0.0 5193	0.5850CCE C3		
5201 52213 52213 52223 52223 52231 52237 52237 52237 52237 52237	C. 750CC0E 02 C. 5C0005C 01 0. 775CCCE 01 C. 1347U0E 01 C. 1347U0E 01 C. 1347C0E 01 C. 1393CCE 02 C. 377CCCE 02		C.13CCCC5 C1 0.6323002 C4 0.437CC3-C1 0.337CC02-C1 0.337CC02-C1 0.337C002-C1 0.337C002-C1 0.337C002-C1	52051 52227 52227 52227 52227 52227 52227 5257 5227 5227 5227 5227 5227 5227 5227 5227 5227 5227 5227 5227 5227 5227 5257 5227 5257 5227 5257 5227 5257 5227 52575 5227 5257 5227 5257 5227 5257 5227 5257 5227 5257 5227 5257 5275757 52757 5275757 5275757 5275757 5275757 5275757 5275757 527575757 5275757575	0.5400000000000000000000000000000000000	C.9CC0CCE 00 52C5 0.285000E 03 5211 C.2C00C0E 00 5217 C.2C00C0E 00 5223 C.2C0020E 00 5229 C.2CCCCCE 0C 5235 0.200000E 00 5241 C.C 5247	0.0 52C6 0.0 5212 C.0 5214 0.C 5224 0.0 52230 C.C 5230 C.C 52342 C.C 5249	C.C C.C C.C C.C C.C C.C C.C C.C C C.C C C C C C C C C C C C C C C C C C C C
5390	0.2700008 03	3351	0.0	5352	C.C 5393	C.585CCCE C3		7
3032	-0.12350CE-05							
8063	-0.1256002-05							
8100	-0.1255000-05							
8132	-0.1255002-05							
9901 9907 9913 9913	0.4000000 03 0.100000000033 0.3000000004 0.20000000000034	5202 5203 5213 5313).1000000 C2).1000000 C3).1000000 C2).3000000 C2	7903 9907 9915 9921	C.2CCOCOE 03 9904 0.5COOCOE 02 9910 C.2COOCOE 03 9916 0.2CCOCCE 03 9922	0.500000E 02 9905 C.200000E 03 9911 C.500000E 02 9917 C.500000E C3 9923	C+C 99C6 0+O 9912 0+C 9518 C+C 9924	0.C 0.0 C.O 0.C

Input 55 Data 1

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	NREG=		5	NRIT=		6 NCAN≕	r	2 N	ZB=	2	KPC=	100	
	1MAX=		12	J.1AX=	1	0 ICA=		3 I	C3=	10	JC I =	2	
	JCE=		٤	DELT=	0.0010	0 SI=	0.50000	DCE 09 W	= (0.100000E	10 V=	C.26000CE 06	
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	έ.Ρ1=	C • C	010	3P2=	0.00	10 DAPF=	• 0•1	1000 0	ELTC=	0.001000	DELTI=	0.001000	
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Core map

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CANA	Lé l									
SEZ.CA4.=	0.130000	P C1)2= (0.7600C0E	02	SC =	C.5400COE 01	R0=	C.800000E	00
PALN1=	2479.	-11=	ડ (0381= C	.286CCOE	С З				
8451		CPT 1			P201		VOL			. •
0.7750 0.2830 0.1347 0.13947 0.6363 0.6363 0.3770	002-01 006-01 006-01 001-01 001-02 002-02].4470].4470].4470].4470].4470].5310	00£-01 00£-01 00£-01 00£-01 00£-01 00£-01		200000 200000 200000 200000 200000 200000 200000 0		0.119383E 0.115383E 0.119383E 0.119383E 0.119383E 0.119383E 0.125324E	02 02 02 02 02 02 02		
CANA	LE 2									
SEZ.CAN.=	0.130000	ē C1	DZ=	C.760000E	02	SC=	0.540CCCE 01	R0=	c.sooocce (oc
PAEN1=	6528.	N 1 =	ن ن	CSE1= C	.236CCOF	C 3				
RESI		CPTI	L		PZG1		VGL			
0.775C 0.2830 0.1847 0.1393 0.6360 0.377C	CCE-C1 CCE-C1 CCE-C1 CCE-C1 CCE-C1 CCE-C2 CCE-C2	0.4470 0.4470 0.4470 0.4470 0.4470 0.4470	CCE-01 CCE-01 CCE-01 CCE-01 CCE-01 CCE-01		•200000E •200000E •200000E •200000E •200000E •200000E	00 00 00 00	0.119383E 0.119383E 0.119383E 0.119383E 0.119383E 0.119383E 0.125324E	02 02 02 02 02 02 02		

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Channel

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REG.	1	ALFA 1	ALFA 2	ALFA 3	ALFA 4
ρ		C•C	C.O	G • C	0.0
REG.)	ALFA 1	ALFA 2	ALFA 3	ALFA 4
P		-0.1255;;;;=05	C•0	0 • C	0.0
REG.	3	ALFA 1	ALFA 2	ALFA 3	ALFA 4
ρ		-0.12553996-05	0.0	0.C	C.O
REG.	ú	ALFA 1	ALFA 2	ALFA 3	ALFA 4
ρ		-0.12353695-05	0.0	0.0	0.0
REG.	5	ALFA 1	ALFA 2	ALFA 3	ALFA 4
ρ		-0.12553998-05	0.0	0 • C	0.0

TEMPERATURE COEFFICIENTS

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REGIONI AVVOLONATO 1 2 3 4 5

ITER42	1001	:	DelT	Ĥ	EΡ	VELENO		IDER
1		0.10	0001-02	0.28735	iD 02	0 • C		2
2		0.10	0C08-02	-0.36366	D 03	0.1999999D	21	1
3		0.10	00012-02	-0.48248	ID 01	0.146459260	20	15
4		C.10	0003-02	0.15345	iD 01	0•12153696D	20	18
5		0.13	0008-02	0.31339	00 00	C.12755075D	20	31
5		0.10	0002-02	0.19258	00-C1	0.129128320	20	31
7		C.1C	CCOE-C2	-0.34035	SC-02	0.129257450	20	
3		0.10	0C0E-02	-0.27554	D-01	0.12938658D	20	
9		0.10	CCCE-02	-0.52522	20-01	0.129515710	20	
10		0.10	CCOE-02	-C.78358	0-C1	0-12964483D	20	
ТЕМРС	TOTALE DI	CALCELO	SEC.	19.090	TEMPO DI	SOLO CRITIC	SEC.	11.340

<u>Criticality</u> <u>Search</u>

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M3.1

3 CATI SU NASTRO

RISULTATI DEL CRITIC

C• :)	0.0) . ()	C • C	0.0	0.0	0.C	0.0	0.0	0.0
0.319059-13	0.51905_ 13	0.526218 13	0.53269E 13	0+337288 13	C.45240E 13	0.258458 13	0.96094E 12	0.37967E 12	0.0
0.242418 14	C+242416 14	J.244338 14	C.24919E 14	0.253788 14	C.21462F 10	C.12267E 14	0.3923CE 13	C.14363E 13	0.0
0.334328 19	C.384321 14	0.839336 14	C.39389E 14	C.905C4E 14	C.75323E 14	0.426748 14	0.137CCE 14	C.5C182E 13	0.C
0.11935: 15	0.115352 15	5.11943C 15	0.119256 15	0.118356 15	0.382398 14	0.55553E 14	0.175318 14	C.65332E 13	0.0
0.119908-15	C.1133CE 13	0.115212 13	0.10310m 15	0.102408 15	0.835038 14	0.477278 14	0.153918 14	C.56441E 13	0.0
0.1007E 15	C.10007_ 13	2,525350 14	C.811336 14	0.72810E 14	0.58357E 14	0.33557E 14	0.10954E 14	0.39884E 13	0.0
0.714998 14	0.714396 14	0.05491E 14	0.553028 14	0.49624E 14	0.392418 14	C.22476E 14	0.72641E 13	0.26690E 13	0.0
0.417719 13	0.317710 14	0.33133E IX	C.327338 14	0.286728 14	C.22575E 14	0.12904E 14	0.41680F 13	0.153C9E 13	0.0
0.105158 14	C.105152 14	0.965128 13	C.8450be 13	C.7540CE 13	0.602348 13	C.34581E 13	0.111118 13	0.40787E 12	0.0
0.22253E 13	0.222636 13	0.200J5E 13	0.130498 13	0.130598 13	0.127428 13	0.72881E 12	0.243598 12	C.10754E 12	0.0
C. 0	C.))•0	G.C	C • C	C.0	0 • C	0.0	0 • C	0.0

FLU(1

Print type m3_

Flux map

These flux maps (pages 60 and 61) correspond to the equilibrium conditions after criticality search. Results are stored on tape. The calculation was restarted from this point. See page 43. i

FLUX 2

0.0	0.0	0.0	0.0	C.O	C.O	0 . C	0•0	0.0	0.0
0.31530E 14	0.31530E 14	0.318878 14	0.32322E 14	C.34250E 14	0.29805E 14	0.1729CE 14	0.68421E 13	0.34206E 13	0.0
0.51329E 14	C.513298 14	0.5136CE 14	0.52;16E 14	0.55098E 14	C.47155E 14	0.27C95E 14	0.114C1E 14	0.62117E 13	0.C
0.145COE 15	C.14500£ 15	0.14536E 15	0.14761E 15	0.15158E 15	C.12722E 15	C.72331E 14	C.3C938E 14	C.17311E 14	0_0
0.19623E 15	0.196236 15	0.19565E 15	0.19558E 15	0.19878E 15	C.16570E 15	0.94035E 14	0.40210E 14	0.22499E 14	0.0
0.20034E 15	0.200348 15	0.19223E 15	C.17314E 15	0.17001E 15	0.14033E 15	0.80621E 14	0.34659E 14	0.19437E 14	O.C Print
0.17072E 15	C.170728 15	0.15767E 15	0.134888 15	0.12002E 15	0.97472E 14	0.56571E 14	0.24437E 14	C.13735E 14	0.0) type nz
0.12201E 15	C.122010 15	0.11142E 15	C.536076 14	C.81797F 14	0.65540E 14	0.37885E 14	0.16354E 14	0.91909E 13	0.0
0.71372F 14	C.71372E 14	C.65003E 14	C.543868 14	C.47318E 14	0.37751E 14	0.21781E 14	0.93954E 13	C.52792E 13	0.0
0.231018 14	0.231013-14	J.21192E 14	0.181136 14	0.16175F 14	C.13124E 14	C.76C81E 13	0.32215E 13	C.17611E 13	0.0
0.13697E 14	0.136878 14	0.12613E 14	0.11100E 14	0.10230F 14	C.33839E 13	C.48623E 13	0.193C6F 13	C.96621E 12	0.0
0.0	C.O	J.U	0.0	0.0	C + O	0 • C	C.O	C.O	0.0
175 = 1	RESIDUAL 0.	41350E 12	RE2 0.52446E	16 TO =	0.0 I	T = 0			
CORE AVERAGE	FLJX	1	FLUX 2	PCWLR	D	P2	REP	TO	Print
	0.45333[1) C.(34948D 14	C.500COE C9	-0.6679	5D 13 -C	•78558D-C1	C • C	ftyp <u>e ne</u>

- 65 -

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CHANNEL	N	1	TIME =	0.100	MASS F	ECW =	0.58500COE	03
3		0.0		0-0	C_ n		0_0	0.0

2	0+0	U. U	C. 0	0.0	0.0	0.0	270.000
4	671.836	567.035	491.000	416.971	342.916	297.564	272.617
5	1037.563	838.768	694.205	553.187	411.859	325.152	277.471
6	1206 • 237	965.557	79 ⁰ •196	618.856	446.878	341.205	283-111
7	1119.143	902-837	745 • 132	59C.818	435.768	340.401	287.982
8	969 .47 6	793.597	665.059	539.177	412.537	334.548	291.690
9	838.803	638.194	595.126	493.920	391.858	328.855	294 • 248
10	643.919	554.410	488•593	423.798	358-298	317.769	295.517

AVERAGE TEMPERATURES

3	0•0	C_C
4	497.551	272.617
5	707.116	277.471
6	805-545	233.111
7	758.752	297.982
8	675.968	291.690
9	603•580	294.248
10	493.804	235.517

<u>Print</u> type no

i

CHAINEL N	2 TIM	E = C.1CO	MASS F	LOW = C.585000	ICE 03]
3	0.0	C•C	C.0	0.0	0.C	0.0	270.000	
4	486.575	430.033	338.902	349.001	309.226	284.952	271.591	
5	716.701	670-754	516.597	434.609	352.552	302.276	274.621	
5	816.607	676.034	573.747	473.795	373.508	311.909	278.042	
7	733.384	613.948	529.930	445.812	361.311	309.347	280.793	
3	620.879	533.035	468.910	406.109	342.953	304.073	282.706	
Э	527.331	454.327	418.219	373.002	327.460	299.380	283.953	Print
10	422.485	336.893	360.793	335.158	309.299	293 • 332	284.561	[[CIM]
								(type
AV ERA DE	TEMPERATURE	S						
3	C - C	c.c						
4	392.823	271.531						
5	524.245	274.521						
ś	532.743	278.042						
7	537.377	230.733						
8	471.331	232.705						
Э	422.063	233.953						
10	352.920	234.531						
T I 4F	0.100	TOTAL	. ENERGY	0.48166E 09	4 V ER A	GE ENERGY	0.544786 CC	Print
26300	INT OF CORE	w I T	TH ENERGY I	LESS THAN				fype_
	100.000		C•.	100J0E 02				
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CORE AVER	R 165	FLUX I	FL	UX 2	PCWER	01	2	REP

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67 I.

TO

n 8

n 9




0.23363 01 0.233536 01 0.219436 01 0.196516 01 0.166726 01 0.113706 01 0.591456 00 0.247016 00 0.761156 01 0.731166 01 0.712436 01 0.636276 01 0.539406 01 0.366076 01 0.189256 01 0.901856 00 0.153318 02 0.158346 02 0.147106 02 0.123296 02 0.108336 02 0.731706 01 0.378216 01 0.460266 01 0.320206 02 0.320206 02 0.227456 02 0.237796 02 0.108436 02 0.130126 02 0.683006 01 0.290296 01 0.572996 02 0.572396 02 0.5583036 02 0.459526 02 0.351536 02 0.236776 02 0.124546 02 0.529286 01 0.115176 03 0.115176 03 0.595536 02 0.753436 02 0.556406 02 0.356116 02 0.183406 02 0.767276 01 0.115936 03 0.115176 03 0.595536 02 0.749716 02 0.556406 02 0.330166 02 0.183406 02 0.767276 01 0.115936 03 0.415376 02 0.469166 02 0.305076 02 0.232066 01 0.237596 01



. 69

n10



RISULTATI DEL CALCOLC AL TEMPO 0.200

IT = 200

This print indicates that at time 0.2 sec. The results were stored on tope. The number (4) indicates that, To restart from this point, the titel card must contain 4 in the Third column and 4 DD statements must be given. (See pages 42, 50, 52)

Some date were changed, they are printed in the new coloulation:

1	0.500000E 01	2	0.20000E-02	3	0.100000E-	• 0 2 4	0.120000E 02	5	0.1000005 02	6	0.500000E 01
3061	0.0	3062	0.200000E 30	3063	0.500000E	00					
3091	0.0	3092	0.245000E 03	3093	0.612000E	03					
9901 9907 9913 9919 9925 9931	0.200000E 03 0.0 0.300000E 03 0.500000E 03 0.600000E 03 0.100000E 03	9902 9908 9914 9920 9926 9932	0.0 0.0 0.500000E 01 0.500000E 02 0.500000E 02 0.250000E 03	9903 9909 9915 9921 9927 9933	0.0 0.0 0.100000E 0.250000E 0.100000E 0.100000E	9904 9910 03 9916 02 9922 03 9928 03 9934	0.0 0.250000E 02 0.100000E 02 0.250000E 02 0.250000E 03	9905 9911 9923 9929 9935	0.0 0.0 0.0 0.0 0.0 0.0	9906 9912 9918 9924 9930 9936	0 • 0 0 • 0 0 • 0 0 • 0 0 • 0 0 • 0

Appendix A

The two diffusion equations are:

$$D_{I} \nabla^{2} \Psi - \Sigma_{R} \Psi + (I - \beta) \cdot (\Sigma_{I} \Psi + \Sigma_{f_{R}} \cdot \Psi) + \sum_{i=1}^{N} \lambda_{i} C_{i} = \frac{1}{W} \frac{d\Psi}{d\xi}$$

- 71 -

2)
$$D_1 \nabla^2 \varphi - \Sigma_a \varphi + \Sigma_R \cdot p \cdot \varphi = \frac{1}{V} \frac{d \varphi}{d z}$$

by reordering and substituting the time derivative with its finite-different expression for a time interval Δt , equation \dot{z}) takes the new expression

$$D_2 \nabla^2 \varphi - A' \varphi + B \varphi + C' = \frac{1}{v} \frac{\varphi - \varphi}{\Delta t}$$

the new expression for eq. d) is formally the same, therefore all the following discussions will refer only to eq. 3) as representative for both fast and thermal group. By putting:

$$A = A' + \frac{4}{v \cdot \Delta t}$$
 and $C = C' + \frac{\varphi^*}{v \cdot \Delta t}$

and reordering, then eq. 3) becomes:

$$D\nabla^2 \varphi - A \cdot \varphi + B \cdot \psi + C = 0$$

This equation is now integrated over the sector



Fig. 5

3)

4)

- 72 D SSS[V1 4] · d2 · nd0 - SSSA4 + SSSB4 + SSSC + 0 5)

For the theorem of Green the volume integral of the Laplace operator of φ , is equal to the integral over the surface of the volume, of the gradient of φ normal to the surface:

$$D \iiint_{\nabla} \nabla^2 \varphi \cdot dz \cdot d\eta \cdot \eta d\theta = \int \int_{\overline{S}} \frac{d\varphi}{d\eta} ds$$

Equation 5) then becomes:

6)
$$-D \iint_{S} \frac{d\varphi}{dn} ds + \iiint A \varphi - \iiint B \psi - \iiint C = C$$

all the signs have been changed in order to have a coefficient matrix with positive entries in the main diagonal and negative all the other entries.

Each integral term of eq. 6) will now be transformed into a finite-different expression where the flux φ_{ij} is considered constant in the volume corresponding to the point (i,j) of the lattice. This volume has been divided into four sub-volumes Vol. 1,2,3,4; within each sub-volume the physical characteristic of the medium are constant.



Here below are the expressions for each term; with the help
of fig. 6 they are self explanatory:

$$D \iint_{S} \frac{k\varphi}{kn} \, dS = - \frac{p_{ij} - q_{i,j+1}}{k_{j}} \left[D_{2} \frac{k_{i}}{2} + D_{3} \frac{\kappa_{i+1}}{2} \right] \cdot \left(2 - \frac{k_{j}}{2}\right) 2 \pi + \frac{q_{i,j+1} - q_{i,j}}{k_{j+1}} \left[D_{4} \frac{\kappa_{i}}{2} + D_{4} \frac{\kappa_{i+1}}{2} \right] \cdot \left(2 + \frac{k_{i+1}}{2}\right) 2 \pi + \frac{q_{i,j} - q_{i,j}}{\kappa_{i}} \left[D_{4} \frac{k_{i}}{2} + D_{4} \frac{\kappa_{i+1}}{2} \right] \cdot \left(2 + \frac{k_{i+1}}{2}\right) 2 \pi + \frac{q_{i,j} - q_{i+1}}{\kappa_{i}} \left[D_{4} \frac{k_{j+1}}{2} \left(2 + \frac{k_{j+1}}{4}\right) + D_{2} \frac{k_{i}}{2} \left(n - \frac{k_{j}}{4}\right) \right] 2 \pi - \frac{q_{i+1,j} - q_{i+1,j}}{\kappa_{i}} \left[D_{3} \frac{k_{j}}{2} \left(n - \frac{k_{j}}{4}\right) + D_{4} \frac{k_{j+4}}{2} \left(2 + \frac{k_{j+4}}{4}\right) \right]^{\frac{1}{2}} \pi$$

1)
$$\iiint A \cdot \varphi \cdot dz \cdot dz \cdot z d\theta =$$

$$q_{ij} A \oplus \left[\frac{\kappa_{ij}}{2} \frac{k_{j+1}}{2} \left(2 + \frac{k_{i+1}}{4}\right)\right] 2\Pi$$

$$+ \varphi_{ij} A \oplus \left[\frac{\kappa_{i}}{2} \frac{k_{i}}{2} \left(2 - \frac{k_{j}}{4}\right)\right] 2\Pi$$

$$+ \varphi_{ij} A \oplus \left[\frac{\kappa_{i+1}}{2} \frac{k_{i}}{2} \left(2 - \frac{k_{j}}{4}\right)\right] 2\Pi$$

$$+ \varphi_{ij} A \oplus \left[\frac{\kappa_{i+1}}{2} \frac{k_{i}}{2} \left(2 - \frac{k_{j}}{4}\right)\right] 2\Pi$$

9)
$$\begin{aligned} & \iiint \mathbf{B} \, \boldsymbol{\psi} = \boldsymbol{\psi}_{ij} \cdot \mathbf{B}_{(j)} \cdot \mathbf{V} \, \boldsymbol{\psi}_{L1} \cdot 2\mathbf{\Pi} \\ & + \boldsymbol{\psi}_{ij} \cdot \mathbf{B}_{(j)} \cdot \mathbf{V} \, \boldsymbol{\psi}_{L2} \cdot 2\mathbf{\Pi} \\ & + \boldsymbol{\psi}_{ij} \cdot \mathbf{B}_{(j)} \cdot \mathbf{V} \, \boldsymbol{\psi}_{L3} \cdot 2\mathbf{\Pi} \\ & + \boldsymbol{\psi}_{ij} \cdot \mathbf{B}_{(j)} \cdot \mathbf{V} \, \boldsymbol{\psi}_{L4} \cdot 2\mathbf{\Pi} \end{aligned}$$

10)
$$\int \int C = C_{ij} V \phi L \ell \cdot l T + C_{ij} V \phi L \ell \cdot l T + C_{ij} V \phi L \ell \cdot l T + C_{ij} V \phi L \ell \cdot l T + C_{ij} V \phi L \ell \cdot l T + C_{ij} V \phi L \ell \cdot l T$$

where VØL 1,2,3,4 . $2 \widetilde{11}$ are the subvolumes of fig. 6, or the expressions in parenthesis $\boxed{}$ of eq. 8),

the integration in θ has been made over the whole circumference and due to the axial symmetry, it is only a multiplication with $2~\pi$. This factor being common to all terms can be eliminated.

The final expression of the equation after reordering is:

11)
$$-\alpha_{ij} \cdot q_{i-i,j} - \beta_{ij} \cdot q_{i,j-1} - \chi_{ij} \cdot q_{i+1,j} - \delta_{ij} \cdot q_{i,j+1} + \epsilon_{ij} \cdot q_{ij} = \delta_{ij} \cdot q_{ij} + c_{ij}$$

where

$$\begin{aligned} d_{ij} &= \frac{D_{i} \frac{h_{j+i}}{4} \left(k + \frac{h_{j+i}}{4}\right) + D_{2} \frac{h_{j}}{2} \left(k - \frac{h_{j}}{4}\right)}{K_{i}} \\ B_{ij} &= \frac{\left(k - \frac{h_{j}}{2}\right) \left[D_{2} \frac{k_{i}}{2} + D_{3} \frac{k_{i+i}}{2}\right]}{h_{j}} \\ d_{ij} &= \frac{D_{4} \frac{h_{i+i}}{2} \left(k + \frac{h_{i+i}}{4}\right) + D_{3} \frac{h_{i}}{2} \left(k - \frac{h_{i}}{4}\right)}{K_{i+i}} \\ d_{ij} &= \frac{D_{4} \frac{h_{i+i}}{2} \left(k + \frac{h_{i+i}}{2}\right) \left[D_{i} \frac{K_{i}}{2} + D_{4} \frac{K_{i+i}}{2}\right]}{h_{j+i}} \\ d_{ij} &= \frac{\left(k + \frac{h_{i+i}}{2}\right) \left[D_{i} \frac{K_{i}}{2} + D_{4} \frac{K_{i+i}}{2}\right]}{h_{j+i}} \\ E_{ij} &= \frac{d_{ij}}{k_{j}} + B_{ij} + Y_{ij} + S_{ij} + H_{2} \left(k - \frac{h_{i}}{4}\right) \left[A_{0} \frac{K_{i}}{2} + A_{0} \frac{K_{i+i}}{2}\right] + \frac{h_{ij}}{2} \left(k - \frac{h_{i}}{4}\right) \left[A_{0} \frac{K_{i}}{2} + A_{3} \frac{K_{i+i}}{2}\right] \\ b_{ij} \psi_{ij} + C_{ij} &= \psi_{ij} \left[B_{i} V_{0} I_{i} + B_{2} V_{0} I_{2} + B_{3} V_{0} I_{3} + B_{4} V_{0} I_{4}\right] + \left[C_{i} V_{0} I_{i} + C_{2} V_{0} I_{2} + C_{3} V_{0} I_{3} + C_{4} V_{0} I_{4}\right] \end{aligned}$$

It can be verified that: $\delta_{ij} = \beta_{i,j+4}$

and Yij = di+1, j

the pentadiagonal matrix of the coefficients is symmetric. The equation 11) can be written:

it is possible therefore to avoid the storage of the two reactor γ and \jmath .

4673C.

Appendix B

Error due to the time discretization in steps of length Δ t.

Let us consider the one-group time dependent diffusion equation:

 $D \nabla \varphi - \Sigma_{n} \varphi + \kappa \Sigma_{n} \varphi = \frac{1}{r} \frac{d \varphi}{d t}$

neglecting the spatial flux distribution and the diffusion term the eq. reduces to:



let $\frac{\delta k}{l} = \alpha$ substituting the time derivative with its finite-difference expression we have:

$$\alpha \varphi = \frac{\varphi - \varphi^*}{\Delta t}$$

at the time $t = n \Delta t$ the value of the flux as determined by this approximation will be :

$$\varphi' = \varphi_{\bullet} \left(\frac{1}{1-\alpha \Delta t}\right)^n$$
 a)

- 76 -

whereas the exact analytical solution will give :

The appropriate solution φ' exceeds the exact one φ'' by an error :

$$\mathcal{E} = \frac{\varphi' - \varphi''}{\varphi''} = \frac{(1 - \frac{\Delta t}{T})^{-n} - e^{m \frac{\Delta t}{T}}}{e^{m \frac{\Delta t}{T}}}$$

In the range of accidents which can be treated by this code, the power level does not become greater than 10³ the nominal value. A scram or a temperature reaction will interrupt the increase of power, and in the great majority of cases the increase of power is less than 10 times the nominal value.

Here below are calculated the values of ξ for various values of the ratio $\Delta \xi / T$ between the time step and the reactor period, for an increase of power of 10³ and 10.

t

Ъ)

 $e^{M\frac{\Delta t}{T}} = 10^{3}$ $m\frac{\Delta t}{T} = 3 \cdot l_{ge} 10 = 6.90776$ $m = 6.9 \cdots / \frac{\Delta t}{T}$ $e^{\frac{1-\Delta t}{T} - \frac{1000}{\Delta t}}$





$\Delta t/_T$	٤
0.1	13.1 %
0.05	6.2%
0.01	1 10

The iterations stop when the residual is smaller than a given value of ξ , criterium of convergence :



 ${\cal E}$ should be smaller than the possible variation of ${\cal q}$ in ${\sf A}{\sf t}$

$$\begin{aligned} \varphi &= \varphi_{o} e^{\frac{\Delta t}{T}} \\ \varepsilon &< \frac{|\Delta \varphi|}{\varphi} = \left| 1 - e^{\frac{\Delta t}{T}} \right| \\ e^{\frac{\Delta t}{T}} &= 1 + \frac{\Delta t}{T} + \int_{v}^{v} \left(\frac{\Delta t}{T} \right)^{2} \frac{1}{2!} + \end{aligned}$$

truncating at the second term and substituting



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Alfred Nobel

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