EUR 4464 e

1

-

EUROPEAN ATOMIC ENERGY COMMUNITY - EURATOM

HEROIC A NEW METHOD FOR CLUSTER-TYPE FUELLED LATTICE CELL CALCULATIONS

by

R. CUNIBERTI and Cl. DAOLIO

1970



Joint Nuclear Research Center Ispra Establishment - Italy

Reactor Physics Department Reactor Theory and Analysis

LEGAL NOTICE

This document was prepared under the sponsorship of the Commission of the European Communities.

Neither the Commission of the European Communities, its contractors nor any person acting on their behalf :

make any warranty or representation, express or implied, with respect to the accuracy, completeness or usefulness of the information contained in this document, or that the use of any information, apparatus, method or process disclosed in this document may not infringe privately owned rights; or

assume any liability with respect to the use of, or for damages resulting from the use of any information, apparatus, method or process disclosed in this document.

This report is on sale at the addresses listed on cover page 4

at the price of FF 9.45 FB 85. DM 6.20 Lit 1,060 Fl. 6.20

When ordering, please quote the EUR number and the title which are indicated on the cover of each report.

> Printed by L. Vanmelle S.A. Ghent, March 1970

This document was reproduced on the basis of the best available copy.

HEROIC - A NEW METHOD FOR CLUSTER-TYPE FUELLED LATTICE CELL CALCULATIONS, by R. CUNIBERTI and Cl. DAOLIO

European Atomic Energy Community — EURATOM Joint Nuclear Research Center — Ispra Establishment (Italy) Reactor Physics Department — Reactor Theory and Analysis Luxembourg, March 1970 — 60 Pages — 13 Figures — FB 85

The HEROIC code described in this report is a cell code which solves the integral transport equation in a 39 groups structure, using the first collision probability formalism.

Collision probabilities are calculated by means of the CLUCOP subroutine developed by Carlvik which deals with a two dimensional cluster geometry.

EUR 4464 e

HEROIC — A NEW METHOD FOR CLUSTER-TYPE FUELLED LATTICE CELL CALCULATIONS, by R. CUNIBERTI and Cl. DAOLIO

European Atomic Energy Community — EURATOM Joint Nuclear Research Center — Ispra Establishment (Italy) Reactor Physics Department — Reactor Theory and Analysis Luxembourg, March 1970 — 60 Pages — 13 Figures — FB 85

The HEROIC code described in this report is a cell code which solves the integral transport equation in a 39 groups structure, using the first collision probability formalism.

Collision probabilities are calculated by means of the CLUCOP subroutine developed by Carlvik which deals with a two dimensional cluster geometry.

e de la companya de l

EUR 4464 e

HEROIC — A NEW METHOD FOR CLUSTER-TYPE FUELLED LATTICE CELL CALCULATIONS, by R. CUNIBERTI and Cl. DAOLIO

European Atomic Energy Community — EURATOM Joint Nuclear Research Center — Ispra Establishment (Italy) Reactor Physics Department — Reactor Theory and Analysis Luxembourg, March 1970 — 60 Pages — 13 Figures — FB 85

The HEROIC code described in this report is a cell code which solves the integral transport equation in a 39 groups structure, using the first collision probability formalism.

Collision probabilities are calculated by means of the CLUCOP subroutine developed by Carlvik which deals with a two dimensional cluster geometry. The code is substantially based on the PINOCCHIO scheme but it is not limited to the description of a restricted number of fixed geometries, and clusters can be treated without any homogenization. The machine time on the IBM-360 is typically of the order of 10 minutes.

The results of the code are compared with fine structure measurements performed at Ispra and elsewhere and it is shown that thermal activation distributions are well reproduced.

The code is substantially based on the PINOCCHIO scheme but it is not limited to the description of a restricted number of fixed geometries, and clusters can be treated without any homogenization. The machine time on the IBM-360 is typically of the order of 10 minutes.

The results of the code are compared with fine structure measurements performed at Ispra and elsewhere and it is shown that thermal activation distributions are well reproduced.

The code is substantially based on the PINOCCHIO scheme but it is not limited to the description of a restricted number of fixed geometries, and clusters can be treated without any homogenization. The machine time on the IBM-360 is typically of the order of 10 minutes.

The results of the code are compared with fine structure measurements performed at Ispra and elsewhere and it is shown that thermal activation distributions are well reproduced.

EUR 4464 e

EUROPEAN ATOMIC ENERGY COMMUNITY - EURATOM

HEROIC

A NEW METHOD FOR CLUSTER-TYPE FUELLED LATTICE CELL CALCULATIONS

by

R. CUNIBERTI and Cl. DAOLIO

1970



Joint Nuclear Research Center Ispra Establishment - Italy

Reactor Physics Department Reactor Theory and Analysis

ABSTRACT

The HEROIC code described in this report is a cell code which solves the integral transport equation in a 39 groups structure, using the first collision probability formalism.

the first collision probability formalism. Collision probabilities are calculated by means of the CLUCOP subroutine developed by Carlvik which deals with a two dimensional cluster geometry.

The code is substantially based on the PINOCCHIO scheme but it is not limited to the description of a restricted number of fixed geometries, and clusters can be treated without any homogenization. The machine time on the IBM-360 is typically of the order of 10 minutes.

The results of the code are compared with fine structure measurements performed at Ispra and elsewhere and it is shown that thermal activation distributions are well reproduced.

KEYWORDS

CODES INTEGRALS EQUATIONS COLLISIONS CLUSTERS GEOMETRY MEASUREMENTS ACTIVATION DISTRIBUTION

HEROIC

A NEW METHOD FOR CLUSTER-TYPE FUELLED LATTICE CELL CALCULATIONS*)

INTRODUCTION

The fuel elements used in most of the heavy water moderated power reactors are characterized by a markedly heterogeneous structure which requires a very refined description of the spatial distribution of neutron events.

Since 1962 an important program at Ispra has been devoted to the assessment of calculational methods for solving the eigenvalue problem in such a reactor cell. The results of this work have been reported by Amyot et al.(1).

Several calculational methods have been set-up and checked against experimental results. Two lattice calculation codes PLUTHARCO(2) and PINOCCHIO(3) are being used and continuously developed.

The first code PLUTHARCO, intended for survey studies, is based on correlation techniques and a phenomenological description. With the second code, PINOCCHIO, an attempt has been made to set-up a model which gives a detailed description of the neutron physics in the cell without having to make adjustments in the broad group data in order to achieve the desired agreement.

Essentially, PINOCCHIO solves the eigenvalue problem for the infinite lattice in a multigroup energy structure. A limited number of well defined cluster geometries can be treated by the code, which uses the first flight collision probabilities approach in the integral transport equation applied to the cell

^{*)} Manuscript received on 18 September 1969.

homogenized in annular regions with a cylindrical outer boundary. The HEROIC code, which is described in this report, is substantially based on the PINOCCHIO scheme, but it uses a generalized cylindrical cell geometry treatment so that clusters can be treated without any homogenization. Furthermore, the code is no longer limited to the description of a restricted number of fixed cluster geometries. Radial and azimuthal Subdivisions of both the pins and the coolant are allowed.

Collision probabilities are calculated by means of the CLUCOP subroutine developed by Carlvik (4) which deals with a twodimensional cluster geometry.

Due to the fact that in the thermal region it is not obvious how the cluster geometry can be simplified and the fact that homogenization always introduces additional sources of errors, the main interest of the HEROIC calculational scheme is in the prediction of thermal events, and, consequently, it can be used as a reference for the test of homogenization procedures. The results of the code are compared with fine structure measurements performed at Ispra and elsewhere, and it is shown that thermal activation distributions are well reproduced.

1. GENERAL

The HEROIC code is a multi-energy group program solving the eigenvalue problem in a reactor cell with a cylindrical outer boundary. The rest of the cell geometry is generally cylindrical so that the cluster geometry can be treated without any intervening homogenization. The program is subdivided in four major parts. In the first part, the calculation of macroscopic cross sections

- 4 -

from input and library data is done. This part includes the evaluation of Dancoff coefficients and resonance integrals. The energy group structure is fixed, and the limits of the groups are given in table 1.

In the second part, the collision probability matrix is set-up. The eigenvalue equations are solved in the third part, producing the flux distribution in space and energy. In the final part, this flux distribution is used to calculate region-averaged macroscopic reaction rates in each group and in three standard sets of groups: "fast", "epithermal" and "thermal". Macroscopic reaction rates for fissile and fertile isotopes and for individual isotopes specified by the user are also calculated. Intensive and extensive parameters of the cell, few group parameters for the homogenized cell, and all the parameters which are generally compared with experiments are calculated in this last section.

2. DATA LIBRARY AND RESONANCE CROSS SECTIONS

The library of HEROIC is exactly the same library available to the PINOCCHIO code. Cell cross section data, except for the resonance integrals, are taken from the GAM-II (5) library for the fast and epithermal region and the GATHER-II (5) library for the thermal range. However, the 2200 m/s constants for U-233, U-235, Pu-239 and Pu-241 have been extracted from the 1965 study of Westcott (6). To construct the HEROIC library in the group structure described in table 1, an energy condensation was performed with the help of GGC-II(5) (for an ideal homogeneous system consisting of D_0O and natural uranium).

- 5 -

Table	1

Table	1
Energy group	structure

Group	High energy	Low energy
n.	(eV)	(eV)
1	1 49 107	1 35 106
2	1 35 10	1 11 10 ⁵
3	1, 11, 10 ⁵	3, 18, 10 ⁴
4	3.18.10	961
5	961	130
6	130	47.8
7	47.8	29
8	29	10.7
9	10.7	2.38
10	2.38	1.5
11	1.5	1.3
12	1.3	1.2
13	1.2	1.15
14	1.15	1.11
15	1.11	1.09
16	1.09	1.07
27	1.07	1.05
28	1.05	1.025
19	1.025	1
20		0.97
21	0.97	0.91
22	0.91	0.7
23	0.7	0.5
24	0.5	0.35
25	0.35	0.31
20	0.31	0.27
28	0.27	0.22
29	0.22	0.16
30	0.16	0.1
31	0.1	0.085
32	0.085	0.06
33	0.06	0.04
34	0.04	0.03
35	0.03	0.025
36	0.025	0.015
37	0.015	0.01
38	0.01	0.005
3 9	0.005	0.0

Diagonal transport corrections are included. Energy exchange between thermal groups in heavy water, light water and graphite is accounted for through the Nelkin - Honeck, the Nelkin and the Parks kernels, respectively. A kernel developed at Ispra (7) describes the interactions with the hydrogen atoms bound in the organic molecule. As in PINOCCHIO, resonance cross sections in a given group are defined simply as being equal to the resonance integrals in the group divided by the lethargy width of the group. Resonance integrals are average values for the cluster obtained by a best fit (8) (9) of the results of the ZUT and TUZ Nordheim's method in the form $I = A + B\sqrt{\text{Seff/M}}$.

In the case of Uranium-235, a slightly more complicated correlation has been developed (10). No attempt has been made to distribute resonance capture in different pins, i.e., inner pins are given a resonance integral equal to that of outer pins.

As suggested by Levine the effective surface is calculated as

$$\frac{S_{eff}}{M} = \frac{1-C}{1+0.1C} \frac{2}{\Gamma P}$$

where g is the fuel density, r the rod radius and C the Dancoff coefficient.

The Dancoff coefficient can be obtained from the fuel to fuel collision probabilities in the limit of black fuel as proposed by Leslie (1;) and is defined as

$$C = \frac{4V_{f}}{S_{f}} / \overline{L_{f}}$$

where V_{f} and C_{f} are the total fuel volume and surface and \overline{L}_{f} is determined by the asymptotic expansion of P(f,f,) for large values of the total macroscopic cross section.

- 7 -

P(f,f) is calculated by the CLUCOP subroutine assuming a total cross-section, $\Sigma_{\rm f} = \frac{1000}{2r}$ for the fuel where r is the radius of a pin.

3. NEUTRON TRANSPORT AND EIGENVALUE PROBLEM

Neutron transport is described by means of first flight collision probabilities. The program is limited to isotropic scattering in the laboratory system (transport corrections are made when preparing the data library).

The method used for calculating the collision probabilities is that described by Carlvik (4) and is included in the CLUCOP code. It is similar to the one used in the PIJ routine (12).

The equations for the method of collision probabilities have been formulated by Amyot et al. in ref. (3) where the eigenvalue problem has been discussed.

The space-energy distribution of the flux for the infinite lattice is calculated with reference to one neutron produced by fission occuring below a given energy limit.

The collision probability equations for the thermal energy groups are solved with the iterative scheme proposed by Honeck (13) (ITER subroutine of the THERMOS code).

The infinite multiplication factor is defined in the framework of a four factor formula, corrected to allow for epithermal fissions

$$K_{00} = \varepsilon \left[p \eta f_{H} + (1-p) \eta f_{api} \right]$$

Cell-averaged cross sections are then produced. The cell averaged diffusion coefficients are calculated by the Benoist (14)formula. Assuming that the buckling remains constant as a function of energy, the criticality equations are solved in the form

- 8 -

$$\sum_{m} \alpha_{mn} \phi_{m} = B^{2} \phi_{n}$$

where

$$\begin{aligned} &\Omega_{mn}' = D_n^{-1} C_{mn} \Sigma_m \quad (m \neq n) \\ &\Omega_{nn}' = D_n^{-1} \left[C_{nn} \Sigma_n - \overline{Z}_{nn} \right] \end{aligned}$$

 C_{mn} is the number of neutrons appearing in group n at each collision in group m; Σ_n is the cell-averaged collision cross-section; Σ_{rn} the cell averaged removal cross-section. D_n the cell averaged diffusion coefficient.

Once the material buckling and the flux spectrum have been obtained, the Keff of the system can be defined as

$$\kappa_{eff} = \frac{\sum_{n} (v \Sigma_{f})_{n} \phi_{n}}{\sum_{n} (\Sigma_{an} + D_{n}B^{2}) \phi_{n}}$$

where all reaction rates are referred to one neutron existing (or produced in) the system.

4. COMPARISON WITH EXPERIMENTAL RESULTS

4.1. Thermal Fine Structure Measurements

The HEROIC code has been tested on thermal fine structure measurements of the neutron density.

Tables 2 and 3 give the results for the seven rod cluster fuel element (carbide fuelled, diphyl cooled) and for the 19 rod cluster ECO fuel element (U metal fuelled and diphyl cooled). Fig. I gives the layout of the 19 rod cluster and subdivisions used in the calculation. Table 4 gives the results for the CIRENE 19 rod cluster (UO₂ fuelled, polystirene cooled D₂O moderated) and table 5 gives the comparison of the fine structure of the thermal neutron density for the Marviken 37 rod cluster (UO₂ fuelled, D₂O cooled and moderated) (Fig. II). For this case the theoretical results obtained by the CLEF Swedish code (18) are also given.

4.2. Neutron Balance Measurements

The HEROIC code has been used for the comparative analysis of substitution and zero-reactivity measurements performed at Ispra, on ECO, and at Montecuccolino, on RB-1, respectively. In fact, the comparison between substitution and zero reactivity techniques is not straight-forward because the measured quantity is the material buckling in the former type of experiment and in K-inf in the latter type.

In table 6 the results of such a comparison in terms of K-inf are given with leakage calculations performed by the two lattices codes PINOCCHIO and HEROIC.

From the comparison given in table 6 it appears that, when a particular accuracy is needed for the interpretation of neutron balance measurements, the more refined code HEROIC gives very satisfactory results.

Fine structure of thermal neutron densities Comparison HEROIC-experiments (fuel element UC/7/25.2 Diphyl)

	ISPRA(15) experimental results	Canadian(16) experimental results	HEROIC
Central pin	0.6738	0.691	0.699
Central can Exter. pins	0.7883 1.0544	1.052	0.813
Exter. cans	1.2618	-	1.288

Table 3

Fine structure of thermal neutron densities Comparison HEROIC-experiments (fuel element U/19/12 Diphyl)

Region	ISPRA(17) experimental results	HEROIC
A	1.	1.
B	1.163	1.176
C	1.792	1.759
1A	2.686	2.620
2	1.561	1.581
3	1.585	1.581
4	1.132	1.159

Fig. I - Details of the detector arrangement in the ECO clustered fuel element for the average flux measurements (dimensions in mm)



Fine structure of thermal neutron densities Comparison HEROIC-experiments CIRENE (fuel element U02/19/0.35 D20 99.7% pitch 30.2 cm)

0.666	0,669
0.773	0.775
1.141	1.140
1.778	1.762
1.801	1.839
2.518	2.506
	0.000 0.773 1.141 1.778 1.801 2.518

Fig. II - Marviken boiler test element. Subdivisions used in the calculations



Cu-63 activation distribution for the Marviken boiler test element (37 rod cluster, D_2^0 mod. and cooled, pitch 24.04 cm). Comparison HEROIC-experiments

Region	Experimental results	HEROIC	CLEF (18) results
1	0.781	0.776	0.784
2	-	0.824	0.845
3	0.892	0.861	0.874
4	0.824	0.817	0.328
5	-	0.877	0.879
6 + 7	0.942 0.934		0.943
8	0.928	0.928	0.929
9		0.996	1.006
10+11	1.056	1.055	1.060
12	1.121	1.124 1.119	
13	-	1.212	1.203
14+15	1.229	1.251	1.245
16	1.385	5 1.405 1.393	
Moderator	1.829 1.750 1.		1.780

<u>Table</u> 6

A comparison of leakage calculations performed by HEROIC and PINOCCHIO in terms of K-inf from different experimental techniques

fuel element	Pitch	K-inf	K-inf	ECO
			TINUCCIIIU	IIBRO I C
U/19/12	18.8 23.5 26.7	1.0902 1.1279 1.1334	1.0906 1.1314 1.1412	1.0905 1.1288 1.1373
UC/7/25.2	23.5 26.7	1.0953 1.1090	1.0927 1.1034	1.0939

Part II

CODE DESCRIPTION

5. SPECIFICATION OF THE CELL GEOMETRY

The composition of the materials which constitute the cell is expressed in the HEROIC input in terms of atomic densities (nuclei/barn cm.). A maximum of ten materials can be considered. If one of the materials is the void (atomic density equal zero) the corresponding macroscopic cross sections is defined by the code as equal to $0.9.10^5$.

The following rules should be followed for the subdivision of the cell in regions.

a) The fuel pins include fuel and canning, and may be subdivided by concentric circles and by diameters.

a) The coolant and the moderator may be subdivided by concentric circles and by radial lines between these circles.

c) Region numbers in the pin should start in the innermost annulus.

d) If an annulus is further subdivided by diameters, these diameters should be specified by giving their angles (in positive orientation) with the line through the cell centre and the pin centre. In this case the numbering of the sectors in the annulus starts in that sector which has the smallest (in degrees) specified diameter as its lower boundary (zero included).

e) The regions numbering inside a pin is always immediately precedent at the number of the coolant region in which the centre of the pin is placed.

f) In the coolant and moderator zones, region numbers start at the innermost regions.

g) If an annulus in the coolant or moderator zones is subdivided by radii, the numbering is done in a recurrent sector in positive orientation.

Fig. III gives an example or regions numbering.

6. ISOTOPES AND LIBRARY

A list of the isotopes that are included in the present version of the code library is given in table 7. Each isotope is identified by two identification numbers (ex. 14002 40). The seven numbers (xyzzz cc) which compose the identification are chosen in the following way:

х

- = 0 : microscopic thermal cross-sections do not depend on temperature (non-moderator isotope)
- = 1 : microscopic thermal cross-sections depend on temperature, that is,the isotope is present in the library at several temperatures (moderator isotopes)
- zzz : identifies the isotope
- cc : isotope temperature in (°C).(=27 for the isotopes in ambient temperature). This number is not given in input for each isotope, but it is assigned to each material, as it is clear that all isotopes which belong to the same material are at the same temperature. Special attention must be paid in the control if these isotopes are included in the library at the asked temperature. For x = 0, cc is always assumed, by the code, equal to 27°C.

- 15 -

The number zzz is also used by the code for the identification of the resonance isotopes in the fuel. For these isotopes (defined resonant in the input data) the coefficients for the resonance integrals calculation are read from library (see section III of the table 7). Through the resonance integrals the microscopic scattering and capture epithermal cross sections are obtained.

Therefore, the HEROIC library is divided into 3 sections:

- a) thermal cross sections
- b) fast and epithermal cross sections
- c) resonance integrals coefficients

The total number of isotopes present in the cluster must be ≤ 50 : in the evaluation of this total number, both the identification numbers have to be considered, that is to say for example that (12001 27) and (12001 70) will be considered by the code as two isotopes.

7. FEW-GROUP PARAMETERS

Few-group macroscopic cross sections for use in diffusion calculation and three groups heterogeneous parameters for use in the SOS code are calculated by HEROIC.

The parameters are averaged values for channel (fuel element and tubes), moderator and cell: they are calculated only if the calculation of the three-groups macroscopic cross sections is requested (see Input Data, word AVE, card 2). Broad group limits for few-group parameters calculations must be given in input by following this rule: from high to low energy, for each group, the lower energy limit must be specified in the input data, by giving the ordinal number corresponding to the lower HEROIC fine group in the broad group (see table 1).

8. CODE LANGUAGE AND SYSTEM

HEROIC, written in FORTRAN IV, can be run on a normal IBM 360/65 following the O.S. nominator, in the HASP system, with overlay organisation. The total number of bytes occupied is \div 284.000, including library subroutines and special functions. The compiler used is FORTRAN H (level zero).

9. EXECUTION TIME

The mean time for a typical cell calculation mainly depends on the accuracy asked for the collision probabilities calculation. The input data parameters for this accuracy, are the number of integration lines and angles (directions) (see input data, word ACC.card 2). Besides, for the same number of lines and directions, the execution time is proporzional to the regions number (see input data, word GEO , card 2). Table 8 shows the approximate times of execution for several numbers of lines and angles. As example, the execution time for the sample case is about of 7.3 minutes. In this case the execution time necessary for the collision probabilities calculation (39 groups, 14 regions, 10 lines, 2 angles) is about 5 minutes. The HEROIC code gives, for this, the progressive times, in seconds, at the end of the subroutines INDAT, INLIB, REGION, AREA, GEOM, BOUND, GLUCOP, RISON, GROUPS, BOUND1, COLLIS, MAIN (see flowdiagram).

- 17 -

The execution time is slightly reduced in the cases after the first if the word MAT (see input data) is not changed. In fact in this case, the library tape must be read completely only once to locate the isotopes of interest (see flow-diagram).

10. INPUT DATA SPECIFICATION

The HEROIC INPUT is controlled by directive words defining different data blocks. The directive words are the following:TIT, GEO, MAT, ASS, ACC, ALB, AVE, STA, END. These words are punched starting in card column 1. The data blocks must be given in the same order as the directive words in the list and, besides, for successive cases, is not necessary to repeat all the input data, but only the directive words, with the corresponding data block, that are to change.

10.1. Title Card (TIT)

Card 1	FORMAT (A3)
cols.	<pre>1 - 3 Specify the word TIT, which defines the data block</pre>
Card 2	FORMAT (18A4)

cols. 1 - 72 Alpha-numeric text (not \$ or '). Should always introduce a new data case.

10.2. Geometry Cards (GEO)

<u>Card</u> 1	FORMAT	(A3)
cols.	1 – 3	Specify the word GEO which defines the block
Card 2	FORMAT	(615)
cols.	1 - 5	NRG, the number of regions in the cell.
		NRG ≤ 30.

- 18 -

- cols. 6 10 NDR, the number of different pin-types with regard to radial subdivision (by concentric circles). $1 \le NDR \le 20$.
- cols. 11 15 NDS, the number of different pin-types with
 regard to sectorial subdivision (by diameters).
 Set zero if no subdivision of this kind is wanted
 O≤NDS≤20
- Cols. 16 20 NP, the number of pin-rings. A pin-ring is a number of pins of the same type and nuclear composition which are equally spaced on a circle around the cell centre. Analogous parts in the different pins of pin-ring have the same region numbers. 1 < NP < 20.
- cols. 21 25 NCORA, the number of radii concentric with the cell boundary (subdivision in radial zones of the coolant and moderator zone). 1<NCORA<30
- cols. 26 30 IKYL, the number of annuli in which subdivision by radial lines is wanted. An annulus of this kind may consist of one or several annuli formed by the CORA; (see card 7). 0 < IDYL <20.
- <u>Card</u> 3 FORMAT (8E10.0)

- PIRAI_{ij}, $j=1,...,K_i$, the radii of the pin-type i. The radii must be given in increasing order and the last will be the pin radius or the external radius of the canning, if there is. The radius PIRAI_i(K_i +1) equal to zero or blank indicates the end of the set. $K_i \leq 30$. Card 3 is repeated NDR times (see Card 2)

Card 4 FORMAT (8E10.0)

- VINGI_{ij}, $j=1,\ldots,L_i$, the diameters dividing the pintype i in sectors. The diameters must be given in degrees in positive orientation from the line through the cell centre and the pin centre. The diameter WINGI_i(1, + 1) equal zero or blank indicates the end of the set. $L_i \leq 30$.

Card 4 is repeated NDS times (see Card 2)

Card 5 FORMAT (515)

- cols. 1 5 NPAN_i, the number of pins in the pin-ring i.
- cols. 6 -10 NPIS_i, the type of pin with respect to radial subdivision in the pin-ring i. Specify a number 1 to NDR (see Card 2)
- cols. 11 -15 NPID_i, the type of pin with respect to sectorial subdivision in the pin-ring i. Specify a number 0 to NDS (see Card 2). Zero means no subdivision of this kind (NDS = 0).
- cols. 16 -20 NS;

Value

Meaning

- O The pins of pin-ring i have not the same region numbers as a pin-ring K(K<i),i.e. the ring introduces new region numbers.
- 1 The pins of pin-ring i have the same region numbers as a pin-ring K (K < i).

- 2 The pins of pin-ring i have the same region numbers as a pin-ring K (K \lt i) but reflected in a line through the cell centre and the pin centre.
- cols. 21 25 NF_i, the pin-ring K(K < i) (see NS_i) is the (i-NF_i)th Specify only if NS_i $\neq 0$.

<u>Card</u> 6 FORMAT (3E10.0)

- cols. 1 10 RIRA, the radius of the pin-ring i. Different pin-rings may have the same radius
- cols. 11 20 PIP_i, the position of a pin in the pin-ring i. Specify in degrees from the boundary of the integration sector, which must be a symmetry axis of the cell.

cols. 21 - 30 RHO_i, the fuel density in the pin-ring i (gr/cm³). Cards 5 and 6 are repeated NP times (see Card 2), starting in the cell centre.

<u>Card</u> 7 FORMAT (8E10.0) CORA, i = 1.....NCORA (see Card 2), the radii concentric with the cell boundary, starting in the centre. Avoid exact coincidence with the radii RIRA (see Card 6).

<u>Card</u> 8 FORMAT (315)

Only if IKYL > 0 (see Card 2)

cols. 1 - 5 NDEL_i, the number recurrent sectors of the annulus i. cols. 6 - 10 IDEL_i, the recurrent sector is divided in IDEL_i subsectors. $2 \leq IDEL_i \leq 20$

cols. 11 - 15 ISYM, Value Meaning 1 The regions of the recurrent sector are all different (no symmetry) 2 The sector is symmetric (see Fig. III) FORMAT (3E10.0) Card 9 Only if IKYL>0 (see Card 2) 1 - 10 RADA, the inner radius of the annulus i. cols. 11 - 20 $RADB_i$, the outer radius of the annulus i. $RADA_i$ cols. and RADB, must coincide with the CORA; (see Card 7) 21 - 30 AIN_i , the starting position for a recurrent sector. cols. Specify the lower angular boundary in degrees. FORMAT (8E10.0) Card 10 Only if IKYL>0 (see Card 2) ANGLE j=2..... IDEL (see Card 8), the inner angular boundaries in the recurrent sector of annulus i. Restriction: ANGLE, and the sector boundaries must not intersect the pins but through the pin centre Cards 8, 9 and 10 are repeated for IKYL times (see Card 2) starting in the centre. FORMAT (2E10.0) Card 11 cols. 1 - 10 G, the integration angle (in degrees). 45° for a quadratic cell, 30° for an hexagonal cell

cols. 11 - 20 B2EX, the experimental buckling to be used in the Keff calculation (leakage terms).

10.3 Isotopes and Materials Cards (MAT) FORMAT (A3) Card 1 cols. 1 - 3 Specify the word MAT, which defines the data block FORMAT (1215) Card 2 cols. 1 - 5 NMAT, the number of materials. NMAT < 10 cols. 6 -10 NE, the number of isotopes present in the cell NE < 50 IEM. j=1....NMAT, the temperatures of the materials. IEM, is the second identification number (cc) for the isotopes belong to the material j (see Isotopes and Library) FORMAT (1315) Card 3 1 - 5 ISTBA, the first identification number (xyzzz) cols. for the isotope i (see Isotopes and Library) 6 - 10 IRIS, cols. Value Meaning 0 The isotope do not require a resonance integral calculation 1 The isotope requires a resonance integral calculation cols. 11 -15 ICON, Value Meaning The fission (if there is) and absorption 0 reaction rate for the isotope i are not calculated and printed 1 The fission (if there is) and absorption reaction rate for the isotope i are calculated and printed.

for 232≼ISTBA ≪ 242 the reactions are alway
calculated.
For NOP = 0 (see word ACC, Card 2) put always
ICON = 0.

- ISMAT, j=1....NMA, the material code number (see ASS, j, Card 2); NMA is the number of the materials in which the isotope j is present. ISMAT \leq NMAT
- <u>Card</u> 4 FORMAT (8E10.0) DENS j=1.....NMA, the density (nuclei/barn cm) of the isotope i in the materials in which it is present. NMA is defined in the preceding card.

<u>Card</u> 3 and 4 are repeated for NE times (see Card 2)

10.4 The Assigning of Materials to Space Regions (ASS)

Card	1	FORMAT	(A3)

cols. 1 - 3 Specify the word ASS, which defines the data block

Card 2 FORMAT (12, 13, 1515)

- cols. 1 2 IMA, optional data; this number can be used to identify
 the material but it is not read by the program. In fact
 the nth card of this type corresponds to nth material
 (n=1....NMAT).
- NREG, the region numbers associated with the material code number IMA; 1 indicates the end of the set
 Card 2 is repeated NMAT times (see word MAT, Card 2)

10.5 Accuracy of Integration (ACC)

Card 1 FORMAT (A3)

cols. 1 - 3 Specify the word ACC, which defines the data block

- 24 -

Card 2	FORMAT (915)
cols. 1 - 5	NIG, the types number of integration intervals
	over the cluster (generally = 1) NIG ≤ 5
-	NINT $i = 1$ NIG, the number of integration lines
	drawn for each type of integration interval i.
-	INTG, the number of Gauss-points in each annulus
	of the moderator. INTG ≤ 4
-	NANG, the number of integration angles (directions).
-	NOP, the output indicator
	Value Meaning
	O Reduced output (see flow-diagram)
	1 Complete output
Card 3	FORMAT (5E10.0)
-	BINT, $i = 1NIG$, the outer radial boundary
	for each type of integration interval i (see Card 2).
	\mathtt{BINT}_{i} must coincide with one of the \mathtt{CORA}_{i} (see
	word GEO, Card 7).
10.6 Albedo Value (A	ALB)
Card 1	FORMAT (A3)
cols. 1 - 3	Specify the word ALB, which defines the data block
Card 2	FORMAT (2E10.0)
cols. 1 -10	ALBEDO, the albedo for the cell boundary, i.e.
	the probability that a neutron, which has crossed
	the cell boundary, will return.
cols. 11 -20	EPS, the convergence criterion for the neutron
	density in the thermal flux calculation. If

EPS = 0 or blank the standard option EPS = 10^{-4} is used by the code 10.7 Few-group Parameters (AVE) Card 1 FORMAT (A3) cols. 1 - 3Specify the word AVE, which defines the data block Card 2 FORMAT (1215) cols. 1 - 5 NRC, the number of the outermost region of the channel (fuel element and tubes). Generally NRC = NRG-NMOD. NRG = total regions number (see word GEO, Card 2). NMOD=radial and sectorial regions number in the moderator zone. cols. 6 -10 NRCM, the few-group number wanted. NRCM < 10 NGRP, i=1....NRCM, the few-groups limits (see few-group parameters)

10.8 Start Directive (STA)

Card 1	FORMAT (A3)
cols. 1 - 3	Specify the word STA, which causes the computation
	to start

10.9 End Directive (END)

Card 1	FORMAT (A3)
cols. 1 - 3	Specify the word END, which causes the computation
	to end. This card should be the last data card.

11. CODE DIAGNOSTICS

The HEROIC code gives several diagnostic as output, for some input errors. Some diagnostics are self-explanatory. The list is following:

- a) ILLEGAL DIRECTIVE IN DATA : (Subr. INDAT)A non-existent directive word is given.
- b) PROGRAM ERROR STOP 1000: (Subr. INDAT and INLIB)A moderator isotope asked, is not included in the thermal library or is not included for the temperature asked.
- c) PROGRAM ERROR STOP 1001: (Subr. INLIB) An isotope asked is not included in the fast and epithermal library.
- d) PROGRAM ERROR STOP 1002: (Subr. INLIB) An isotope asked as resonant (IRIS = 1, see word MAT Card 3) is not included in the resonance integrals coefficients library.
- e) PROGRAM ERROR STOP 1003: (Subr. INLIB) The density gives for the U238 (see word MAT Card 4) es equal to zero.
- f) ALL REGIONS ARE NOT ASSIGNED A MATERIAL: (Subr. CLUCOP) A control is necessary for the data block ASS.
- g) NEGATIVE VOLUME NOT PERMITTED: (Subr. CLUCOP) A control is necessary for the data block GEO.
- h) ERROR IN GEOMETRY SPECIFICATION: (Subr. GEOM)As for the preceding diagnostic.
- i) PROGRAM ERROR STOP 1004: (Subr. RISON) As for the STOP 1002.
- 1) PROGRAM ERROR STOP 1005: (Subr. RISON) As for the STOP 1002
- m) **** THERMAL FLUX SPECTRUM CALCULATION HAS NOT CONVERGED. PLEASE TRY WITH A GREATER VALUE OF EPS****(see word ALB Card 2) If some diagnostic is printed out, the calculation is stoped for the present case, but the code starts the successive case, if there is.

12. OUTPUT LIST

If NOP = 1 (see word ACC Card 2) a complete output is given
(about 6000 lines).
If NOP = 0 a reduced output is given (about 1000 lines).

See the flow-diagram for the contents of two output types, which is, in any case, self-explanatory.





Library 1 - STCLØK, IHCFCVTH, IHCFCØMH, IHCUØPT, IHCTRCH, IHCSLØG, IHCSEXP, IHCFRXPR, IHCFRXPI, IHCSSORT, IHCFIØSH, IHCUATBL, IHCFINTH

Library 2 - IHCFEXIT, IHCSSCN , IHCSATN2 , IHCFSLIT , IHCSASCN

Library 3 - BESI# , BESK# , BIO# , BI1# , BKO# , BK1#













- 35 -



- 36 -



	1 2	3	4	5 6	57	8 9	10	11 1	2 13	1415	516	17 18	3192	021	22	2324	25	2627	28	930	31	323	334	35	3637	38	394	041	42	3 4	45	464:	748	<i>(</i> 95	(51	52	5354	555	5657	585	960	51 6	263	646	566	67 6	869	70	717	27:	747	576	77.	787!	980
.t					++	+	+			_			++	╋	┝┼	┿	+		$\left \right $	+	┢╌┤	+	┽╌	$\left \right $	+-	$\left \cdot \right $	+	┢	+	-	+	$\left \right $			┼─	$\left \cdot \right $		┿┥	+	H	┼┨				+	++	┿	┝╋	┿	┿	┢┿	╺┿┙	\mathbb{H}	┿	┥┥
1			L			-							+	+		-		_		-			+	$\left \right $	_		_	╞	╽╽	_						\downarrow	-	╎				_		\square				\square	\downarrow	\downarrow	\square	\downarrow	\square	+	
2	JC	ľ	7	//	2 5	. 2			D	20		/	D	P		<u>)</u> 9	$\left \cdot \right $	73		1	P	AS	S	0	2	0	C	M	 .	4					1	\square			_		\downarrow	4	\downarrow							\downarrow	$\downarrow\downarrow$	-	Ц	\perp	\square
30	E	0	$\left \right $	-	\downarrow	_					+					4-	\square	- +	┞┤	\downarrow		-	+				_		\downarrow						-	┢			_				\square			\square				\perp	\square		Ц	\bot	\square
4		1	1	4	\downarrow	_	1			0				2		1	0			0			-	\square						\downarrow						\square		\square																	
5	I .	2	6					1	. 3	7 5						_			\square				-						\square																										
6				1			1			0				2			0																																						
7) .	0						0	. 0					1	3	. 5	4								ĺ																					[]					\square		Π		Π
8				6			1			0			(2			0						Ι					Γ	Π						Τ	\square											Τ	Π			Π			T	Π
9	2.	9	5					0	. 0					1	3	. 5	4							Π				Γ					T														Τ	Π	T	T	Π				
10	I .	5	4	8 9				З	. 0					4		4				T	4	. 5	5					4	1.	7 5					5	$\left \cdot \right $	0 5					5.	2	0					7	. 2	0				Π
11	€.	2						1	1.	28	3	8		Τ			Π						Τ	Π														T									1	T		1				1	Π
12	3 0).	0			_		0	. 0	00) 3	9 2	2						Π					Π											Τ										-		T		1	1			\square	-	Π
13	4 4	T															Π		Π				Τ	Π					\square																			Π			Π			1	Π
14				4			8			27			2	7		2	7			2 7			T	Π					T										1								-	T	-	1	Ħ			+	
15	1 2	0	0	1			0			0						T			Π	T			Τ		1			T							1		Ť.		-			.						+	+					-	
16	1	7	9	3 2	2	- () 4							1		+				+		1	Ť	tt	1			t		T					1-	+ +	+	<u>+-</u> +	+	†			\uparrow	-+	-		+	Ħ	+	+	++	+	H	+	
17	10	0	1	2			0		Π	0						-		+	††	+		┥		Ħ	1		T	ϯ	+	+	d			+	1-	+	+	++	-+	++	┼╂	+			+-	╏	+		+	+	┝	-+-+	┝╺╋	+	+
18	3.	2	6	20)	- 0) 2						\dagger	t		+-		1	<u> </u>	╈		+	+	$\uparrow \uparrow$				╀	\dagger	-		-+-		+		-+	+	$\left \right $	+		┼╂	+	+		+-	┥┥	+		-	+-	++	++	┢┼┼	+	+
19	T	2	3	5		+	1			1			++	╓		+-		+	╞╌╞	╉		+	+	┟┼	+	H		╀	\dagger	\uparrow	+	+	+	-			+	┝╴┼			┼╏		+		+	┟┼	╈	┼╂	+	+	╞╌╋╸	+	┟╼┿	┿	+
20	2.	3	6	6 2	2	- () 4	H					\dagger	┢	╞┼	+-	$\left \right $	+	$\left \right $	+			+-	╉╋	+	+		╉	╀╂		╀┦		+				+-	┥┽				+	┼┤		+	$\left \right $	+	╈	+	+	┢┼┥	+-	┝╋	+	Η
21		2	3	8			1	H		1	+		╉╉		$\left \right $	+	+	-	$\left \right $	+-	┠┼	+	+	┼┼	-+-	H		╀	+	+	H	$\left \right $	+	+	H	┝╌┞	- +	$\left \right $	+	\vdash			+			$\left \cdot \right $	+-	┼╂	+	+	\mathbb{H}	++	┝╋	+	+
22	3.	2	3	83		- (2 2	$\left \right $	+				╋╋	╋	┝┥	╉	$\left \right $	+	┝┤	+	$\left \right $	+	+	┝┼	+	$\left \right $	+	╀	╂╂	+	+	$\left \right $	+		+	$\left \right $	+	$\left + \right $	+	$\left \right $	┼┨		+			┝╌┿	╇	┼╂	+	+-	╟	+	┢┼╋	+-	+
22	1		1			-	-	$\left \cdot \right $	+	\vdash	+	$\left \right $	+	+	$\left \right $	╋		+	$\left\{ + \right\}$	+		+	+-	$\left \right $	+	H	+	╀	$\left\{ \right\}$	+	+	-+-	+		H	\square	+	$\left \right $	+		┼╏	+		+	+	$\left \cdot \right $	+	╢	+	+	┢┼┥	+	┝╌┼	+-	+
,	+	12		7	5	-	12		3	2 0		$\left \right $		+	$\left \right $	╋		╉	$\left \cdot \right $	+	$\left \right $	+	+	$\left \right $	+	H	-	╂	++	+	$\left \right $		+	+	H		╀		+		┼╂	+-	+			$\left \right $	╇	╢	-+	+-	H	+	\vdash	+	+
25	+	•	2	7			0	ľ	+			┼┼	ſ,	+	┞┤	+	H	╋	H	╉	H	╉	╀	┼┼	-+-	+		╀	┼┼	+	H	+	╉┥		H		+-	$\left + \right $	+		┼╏	+	+		+	$\left \cdot \right $	+	╢	+	+	H	+	$ \rightarrow$	+-	+

15. SAMPLE CASE INPUT (see Fig.IV)

ו 38י ו

[1	2	3	4 5	6	7 8	9 1	211	12	1314	415	1617	718	1920		222	324	252	627	282	930	313	233	334	353	637	738	394	041	42	34	445	464	748	485	d 51	52	35	455	5657	758	5960	61 6	263	646	5366	676	5 86 5	370	71	27.	374	757	767.	778	79E	80
5 5		8	94	8		0 2	$\frac{1}{2}$				+-	+		+					+		$\left \right $							╋		+-		-	+		┢	H					-	\vdash	+	┢┼	+	++	┿	┢╋	Π	┿	┿┥	╞┼	┿	╉┥		-
ţ,	2	0 0	0 2		+	+	+	┢╌┼	-+	0	+	\mathbf{H}	4	+	+	╉┫	-	+	+	+		+-	-		+			┢	$\left \right $	+	+		+	┝╌┾╴		+	-+		╀	+	+			$\left \right $	╉	+	┽	┼╉	\square	+	┽┥	┝┼	+	+	$\left \right $	-
	F	6	2 3		+_			$\left \right $	+		+	+		+	+	+		+	-+	+		+	-	+-	+			╉─	$\left \right $		+		+			$\left \cdot \right $	+	+	+	┼┼	+	\vdash	+-	$\left \right $	╀	++	╉	╢	$ \vdash$	+	┯	┝┼	╋	╉┦		-
F	+		63		-		; -		+	1	+	+	1		+	+	2		+	3		+		4	+-	+	+	╋	┼┼	+	+	+	+			$\left\{ \cdot \right\}$	╉	╉┤	+	╀┨	+	-+	+	┼┼	+-	┼┼	+	┤┨	\square	+	┽┥	┢┼	+	+		-
]	╀╴	-	+	+	+			+	+	$\left\{ \cdot \right\}$	+			+			-	+	- r		1	+-		L.			0		$\left \right $	-+-	+	+	+	\vdash	┢	┥┥		┥┥	-	┥┤		-+	+	┼┼		┼┼	+	┼╂	⊢	+	┦	┝┼	╉	+		-
	·	c	+	++			+	$\left \cdot \right $	+	┼╌┥	+	+		Ĥ	+	+	-	+	-	-	$\left \right $	-	+	\vdash	+			+	$\left\{ \cdot \right\}$	-+-	+-	+	+		-	+	+	+	-	┼┤	-+-		┿	┼┼		+	+	┼╂	<u> </u>	+	+	┼┼	╉	╉┦	$\left \right $	-
F	-	-	+		+	H		┼┼		$\left \cdot \right $		+	+	┦┤	-+-	+				+	-+-	+	+	$\left \cdot \right $	╀	+	-+-	╋	$\left\{ \cdot \right\}$	+	+	-	+-		┨─	$\left \right $	-	$\left \right $	+	╄┥		-+		┼┼		+	+	┼┨	\vdash	+	+	┼┼	+	╇┥	$\left \right $	-
		-+			+	$\left \right $	<u>'</u>	$\left\{ \cdot \right\}$	-+-		+	$\left \right $	•	₽	\square			+	-	+	$ \cdot $	+	\parallel		+			+	$\left \right $	+	\square		+		-	+	-	┿	-+-	┼╌┥		+		$\left \right $	+	+	+	┼┨	\vdash	+	\downarrow	┼┼	+	+		
ľ							1	$\left \right $	+	7	+-	H		\mathbb{H}	+	$\overline{+}$	-	+	-	+	\vdash	╞	+					+	$\left \right $	+	+				+	$\left \cdot \right $	-	$\left\{ \cdot \right\}$	+	+		+		$\left \right $		$\left\{ \right\}$	+	┼┨		\vdash	$\downarrow \downarrow$		_	+	$\left \cdot \right $	_
	-	\rightarrow		<u>'</u>			+	$\left \right $	+	\downarrow	_		_	\mathbb{H}		+		+			\vdash	_			+			-	$\left \right $	+	\square	+	_			┝╌┥		┝┤	+	$\left \cdot \right $	-			$\left \right $	-	$\left \right $	+	┿┨	┝╼┥	┢┼┥	\downarrow		+	+	$\left \right $	_
				\$	+		<u> </u>		-	7			8	1		1	2	+		1 3				4	-+-		-	4	$\left \right $		+				-	\downarrow	-			\downarrow	_			\downarrow	_	╁┼	+	.+.4	┢──┥			┞╌┽	-	\square	\mid	_
		C	-		+		1	\downarrow				\downarrow		\downarrow				\perp				4-		\square	-						4	4				$\left \right $		$\left \right $	_	┢╶╁			_	$\downarrow\downarrow$		\downarrow	\downarrow			┝╺┷	!					
7	+		1		_		1			2	_		2	<u>-</u>			0					_	\downarrow		-+-					_	•+		-							$\downarrow \downarrow$	-		_	_↓		┦╢	4	44		⊢∔	-	\downarrow	_			
34		5	5			\square		\downarrow	_	\square						-						1			1					_			1						\perp							\square			Ц	\square		\square				_
P	L	B	_								_			Ľ											-			1-							.			\square						Ц			_									
							0	$\left \cdot \right $																																																
4	V	Ε																																																Ιİ						
2			1 1	ŀ			2			9			39	,																															T					Í						
3	T	A																Τ											ĪĪ						Ι											T		Ť		- +	T	Π		\square		
4	EN	D	T	Π			T					Π		Π				Τ					Π		1	Π		1		T					Ť	f-+			T	\square			T				-				\square	H		\square		_
			T				T		╡	$\uparrow \uparrow$		Π		T				1					Π		1			1		T					1-				1	11	-				\uparrow	$\uparrow \uparrow$	1				$\uparrow \uparrow$	Π				
ſ	T		1			Ħ	T		+		1	\dagger		11		╈		+		\mathbf{T}		┢	\uparrow		1	Н	\vdash	t	$\uparrow \uparrow$	-		+	+		1	<u>†</u> †	╋	$\uparrow \uparrow$	+				\uparrow	\dagger	+	\uparrow	+				$\uparrow \uparrow$	$ \uparrow $	\uparrow	\ddagger	$\left \right $	
F	T		╡		\uparrow	$\uparrow\uparrow$	t	\uparrow	+		+	\square		11		\uparrow		┢		1		+			\uparrow	\dagger		╀	\dagger	-		+			\uparrow		╈		+-	$\uparrow \uparrow$	-		╈	\dagger	1		+			-+	\dagger	$ \uparrow$	\uparrow	\uparrow	$\left \right $	
ŀ	\uparrow		╈		\uparrow	++	╋		+	+	+-			\mathbf{H}		+	┝╼╋	+	$\left \right $			+	+		-	╉┥	$\left \right $	╋	\dagger	+		+	+		1-	╞╶┼╴	+	$\uparrow \uparrow$		╞┼╿		$\left \right $	+	┼╌┼		\uparrow	+	+		i t	+	╞┼	+	+	┝╺┽	
F	1		╈	+	+	++	\dagger	+	╉		+	+		11		+	$\left \right $	+-				+		$\left \right $	-	H		╋	\dagger	+	$\uparrow \uparrow$	╈			\uparrow	┝╍╋	-+	$\uparrow \uparrow$	+				+		+	┼┼	+	┼╉		í†	+ +	╞┼	+	+		-
ł	†-	H	\uparrow		1	╞┼	╋	\uparrow	+		-	\square	1	\mathbf{H}				╋	+				+	+	╈	\mathbf{f}		1-	$\uparrow \uparrow$	╉	+	+		+		┝╶┼╴	+-	\dagger	\uparrow	\dagger			+		+	$\uparrow \uparrow$	+	┼╉		-+-	+	$ \uparrow$	+	$\uparrow \uparrow$	╞┼	-

2851 SYS1.CMPLIB 2851 VOL SER NOSE EURSYO	PASSED	16.SA
2851 SYSOUT	SYSOUT	
1851 SYS1.FORTLIB	KEPT	
2851 SYS1.LIB	KEPT	
2851 VUL SER NUS= EURSY2. 2851 SYS1.SSPLIB	KEPT	
2851 VOL SER NOS= EURSY2. 2851 SYS69101.1125507.RP001.PINOCCHI.R0000290	DELETED	
2851 SYS69101.1125507.RP001.PINOCCHI.GOSET	PASSED	
2851 VUL SEK NUS= EURSY3. 2851 SYS69101.1125507.RP001.PINOCCHI.GOSET	PASSED	
LKED DATE=69.04.11 BEG.T.=17.024 DURATION=0. DESEC PGM=*.+STEL.LKED.SYSLMOD,COND=(+GL,LT,+): [05F001 DD DDNAME=SYSIN [06F001 DD SYSUIT=4	•007 N.OPER=463 STEL.LKED)	00032120 00032130
O7FOOL DD UNIT=SYSCP, LABEL=(,NL), DISP=(MOD, PASS),	SNAME=PUNCH1,	C00032150
(SUDUMP DD SYSOUT=A <u>1.FT09F001</u> DD <u>DSNAME=LIRE</u> ,UNIT=L91, <u>DISP=(OLD,DELETE</u> <u>VOLUME=(PRIVATE.SER=LIBPIN</u>),), <u>LABEL=(1,SL)</u> ,	00032170 C
DCB=(,RECFM=VB,LRECL=800,BLKSIZE=804,DE/ DDDSNAME=+LIBPIN;UNIT=2314,SPACE=(CYL,()	N=2) 2,1)),	с
$\frac{1.FT11F001}{100} DD DSNAME = +4VSIND, UNIT=2314, SPACE=(CYL, (1))$	5,11)),	С
D • SYSIN DD * 236I ALLDC • FOR PINOCCHI GO B 237I JOBLIB DN 1CO 237I PGM=*•DD ON 232 237I FT05F001 DN 00D 237I FT09F001 DN 333 237I FT10F001 DN 230		
2371 FTIIFÖÖI ÖN 231		

- 40 -

***** HERDIC HETEROGENEOUS REACTOR ON INITIAL COLLISION CALCULATION ***** ================ D20 / D20 99.73 / PASSO 20 CM. 1/25.2 1BER OF REGIONS = 14JIAL SUB-DIVISIONS OF RODS TYPE 1 RADII 1.26000 1.26000 1.37500) CIRCLE RADIUS START POS. NR OF RODS ROD TYPES NS NF RHO 0.0 0 13.54000 1 0 0 2.95000 2 0.0 1 0 Ó 0 13,54000 6 3-DIVISIONS (ANNULI) IN CODLANT ZONE AND MODERATOR ZONE 11 1.54890 3.00000 4.40000 4.55000 4.75000 5.05000 5.20000 7.20000 9.20000 11.28380 **FEGRATION SECTOR =** 30.00 PERIMENTAL BUCKLING = 3.92005-04 **APERATURES AND ATOMIC DENSITIES OF PURE MATERIALS** TERIAL 2 1 3 4 27 1.7932E-04 1.(C) 27 27 27)Ŏ1 0.0 0.0 0.0 3.2620E-02 2.3662E-04)12 235 238 0.0 0.0 0.0 0.0 0.0 0.0 3.2383E-02 0.0 0.0 0.0 ĴĨĞ 3.3876E-03 5.8948E-02 0.0 3.3208E-02 0.0 27 0.0 0.0 0.0 6.6236E-02 0.0 0.0 0.0 1.0000E-09 1.0000E-09 1.0000E-09 1.0000E-09 **TPUT INDICATOR(COLL) = 0** CURACY OF INTEGRATION INTERVAL 0.0 4.55000 10 INTEGRATION LINES **ABER OF GAUSSPOINTS =** 2 **MBER OF INTEGRATION ANGLES** 2 3EDO BETA =1.00000 JTRON DENSITY CONVERGENCE CRITERION FOR THERMAL FLUX CALCULATION = 0.00010 ACTION RATES ARE REQUESTED FOR FOLLOWING COMBINATIONS OF REGIONS 4BINATION REG. 1-11 REG. 12-14 12 4BINATION V-GROUP PARAMETERS ARE REQUESTED FOR FOLLOWING COMBINATIONS AND GROUPS 1 2 **IBINATION** GROUPS 1- 9 1BINATION GRÖUPS 10-39 SRESSIVE TIME(SEC.) AT THE END OF INDAT = 1.480

GRESSIVE TIME(SEC.) AT THE END OF REGION = 22.840 GRESSIVE TIME(SEC.) AT THE END OF AREA = 22.840 GRESSIVE TIME(SEC.) AT THE END OF GEOM = 27.300 GRESSIVE TIME(SEC.) AT THE END OF BOUND = 27.320 NCOFF COEFFICIENT FOR THE FUEL IN THE CLUSTER Ξ 3.9092E-01 UARE ROOT S/M = 2.6214E-01GRESSIVE TIME(SEC.) AT THE END OF RISON = 27.640 GRESSIVE TIME(SEC.) AT THE END OF GROUPS = 27.740 GRESSIVE TIME(SEC.) AT THE END OF GLUCOP = 27.760 GRESSIVE TIME(SEC.) AT THE END OF BOUND1 = 312.800 GRESSIVE TIME(SEC.) AT THE END OF COLLIS = 313.180 FFUSION CDEFFICIENTS(CELL-AVERAGED)

2.3607E 00 1.1690E 00 1.1873F 00 1.1873F 00 1.2454F 00 1.2335E 00 1.2328E 00 1.22308E 00 1.22306E 00 1.2296F 00 1.2296F 00 1.1721E 00 1.1751E 00 **1.1**715E QO 1.1715E 00 1.1729E 00 1.1729E 00 1.1719E 00 1.1722E 00 1.1722E 00 1.1701E 00 1.1681F 00 1.1652E 00 1.1603E 00 1.1507E 00 1.1442E 00 1.1362E 00 1.1281F 00 1.1128E 00 1.0847E 00 1.0468E 00 9.9539E-01 9-5530E-01 9.2458E-01 8.9146E-01 8.2436E-01 7.6871Ē-0Ī 7.0512E-01 6.3264E-01

5.6655E-01 4.6005E-01

JX SPECTRUM(FLUX*VOLUME)

)UP 1

>138E-01 5.0394E-02 8.5985E-02 1.6171E 00 2.4617E-01 1.9419E-01 4.6372E-01 1.1686E-01 1.5226E-01 2.2257E-01 .058E-01 1.2923E 00 1.1648E 00 1.3281E 00

787E-01 1.0043E-01 1.6765E-01 2.9795E 00 4.8753E-01 3.8792E-01 9.9686E-01 2.6003E-01 3.4406E-01 5.1352E-01 3071E-01 3.2905E 00 3.1423E 00 3.5252E 00 NUP 3

5899E-01 3.1969E-02 5.3865E-02 9.1167E-01 1.6959E-01 1.2912E-01 3.9013E-01 1.0978E-01 1.4840E-01 2.2830E-01 1781E-01 1.7552E 00 1.9930E 00 2.3925E 00 DUP 4

 3036E-01
 5.5624E-02
 9.4749E-02
 1.7327E
 00
 3.2816E-01
 2.2999E-01
 8.1774E-01
 2.4861E-01
 3.4542E-01
 5.4331E-01

 3429E-01
 4.5002E
 00
 5.7431E
 00
 7.3210E
 00

 DUP
 5

.

43

1

+143E-01 2.7438E-02 4.6135E-02 8.6952E-01 1.6959E-03 1.1235E-01 4.1730E-01 1.2899E-01 1.7916E-01 2.8388E-01 +869E-01 2.4202E 00 3.2044E 00 4.1573E 00 DUP 6

107E-02 1.2535E-02 2.1085E-02 4.0160E-01 7.9265E-02 5.1588E-02 1.9707E-01 6.1563E-02 8.5663E-02 1.3605E-01 1335E-02 1.1722E 00 1.5743E 00 2.0564E 00 DUP 7

i694E-02 **5.1971E**-03 **8.7809E**-03 **1.6944E**-01 **3.5018E**-02 **2.1882E**-02 **9.0212E**-02 **2.8817E**-02 **4.0274E**-02 **6.4260E**-02 **3781E**-02 **5.6542E**-01 **7.7196E**-01 **1.0138E** 00 **DUP 8**

2520E-02 1.0472E-02 1.7623E-02 3.4342E-01 6.9737E-02 4.3605E-02 1.7725E-01 5.6485E-02 7.8877E-02 1.2580E-01 5102E-02 1.1045E 00 1.5131E 00 1.9923E 00 DUP 9

L859E-02 1.4397E-02 2.4226E-02 4.7625E-01 9.7526E-02 6.0124E-02 2.4955E-01 8.0063E-02 1.1193E-01 1.7877E-01 3991E-02 1.5795E 00 2.1814F 00 2.8830E 00 JUP 10

1871E-02 4.6190E-03 7.7320E-03 1.5055E-01 2.9435E-02 1.8832E-02 7.2536E-02 2.2795E-02 3.1735E-02 5.0469E-02 466E-02 4.3761E-01 5.9825E-01 7.8927E-01 1UP 11

750E-03 1.8015E-03 3.0126E-03 5.8245E-02 1.1290E-02 7.3063E-03 2.7553E-02 8.5984E-03 1.1951E-02 1.8984E-02 1.6365E-01 2.2268E-01 2.9339E-01 UP 12

1099E-03 9.6740E-04 1.6226E-03 3.1286E-02 6.0985E-03 3.9474E-03 1.5032E-02 4.7055E-03 6.5471E-03 1.0402E-02 +527E-03 8.9849E-02 1.2234E-01 1.6119E-01 IUP 13

195E-03 4.9236E-04 8.2772E-04 1.5955E-02 3.1628F-03 2.0263E-03 7.8940E-03 2.4867F-03 3.4638E-03 5.5121E-03 1916E-03 4.7961E-02 6.5554E-02 8.6440E-02 10P 14

782E-03 4.2623E-04 7.1539E-04 1.3803E-02 2.7438E-03 1.7515E-03 6.8254E-03 2.1494E-03 2.9927E-03 4.7643E-03 991E-03 4.1500E-02 5.6745E-02 7.4830E-02 7.4830E-02 JUP 15

-44

t

188E-03 2.1904E-04 3.6780E-04 7.0982E-03 1.4110E-03 9.0085E-04 3.5165E-03 1.1082E-03 1.5435E-03 2.4570E-03 1889E-03 2.1401E-02 2.9266E-02 3.8595E-02 1UP 16

376E-03 2.2274E-04 3.7407E-04 7.2206E-03 1.4352E-03 9.1535E-04 3.5793E-03 1.1283E-03 1.5717E-03 2.5019E-03 1.25E-03 2.1790E-02 2.9801E-02 3.9301E-02 1.4352E-03 9.1535E-04 3.5793E-03 1.1283E-03 1.5717E-03 2.5019E-03 1.025E-03 1.25E-03 1.1283E-03 1.5717E-03 2.5019E-03 1.25E-03 1.25E-03 1.5717E-03 2.5019E-03 1.5717E-03 1

579E-03 2.2672E-04 3.8084E-04 7.3524E-03 1.4607E-03 9.3303E-04 3.6460E-03 1.1496E-03 1.6017E-03 2.5494E-03 1.75E-03 2.2199E-02 3.0360E-02 4.0038E-02 1.4607E-03 9.3303E-04 3.6460E-03 1.1496E-03 1.6017E-03 2.5494E-03 1.0017E-03 1.001

386E-03 3.0065E-04 5.0425E-04 9.7359E-03 1.9303E-03 1.2335E-03 4.7927E-03 1.5083E-03 2.1001E-03 3.3423E-03 5.2393E-02 5.2393E-02 5.2393E-02 1.9303E-03 1.2335E-03 4.7927E-03 1.5083E-03 2.1001E-03 3.3423E-03 1.92035E-03 1.9205E-03 5E-0

769E-03 3.0791E-04 5.1623E-04 9.9690E-03 1.9744E-03 1.2623E-03 4.8955E-03 1.5399E-03 2.1438E-03 3.4116E-03 893E-03 2.9667E-02 4.0533E-02 5.3443E-02 NUP 20 405E-03 3.7863E-04 6.3485E-04 1.2260E-02 2.4248E-03 1.5518E-03 6.0116E-03 1.8905E-03 2.6321E-03 4.1879E-03 965E-03 3.6397E-02 4.9716E-02 6.5547E-02

909E-03 7.7930E-04 1.3081E-03 2.5259E-02 4.9965E-03 3.2000E-03 1.2429E-02 3.9124E-03 5.4496E-03 8.6703E-03 481E-03 7.5359E-02 1.0296E-01 1.3575E-01

670E-02 3.4481E-03 5.7839E-03 1.1160E-01 2.2055E-02 1.4139E-02 5.4685E-02 1.7188E-02 2.3929E-02 3.8067E-02 964E-02 3.3062E-01 4.5131E-01 5.9489E-01 UP 23

370E-02 4.3793E-03 7.3496E-03 1.4176E-01 2.8129E-02 1.7995E-02 6.9930E-02 2.2006E-02 3.0644E-02 4.8766E-02 579E-02 4.2420E-01 5.7941E-01 7.6379E-01

905E-02 2.5402E-03 4.2662E-03 8.2383E-02 1.6465E-02 1.0475E-02 4.1146E-02 1.2988E-02 1.8096E-02 2.8818E-02 121E-02 2.5147E-01 3.4413E-01 4.5383E-01

222E-02 2.2293E-03 3.7484E-03 7.2610E-02 1.4695E-02 9.2476E-03 3.7052E-02 1.1762E-02 1.6406E-02 2.6159E-02 735E-02 2.2961E-01 3.1543E-01 4.1641E-01 IUP 26

- 45

049E-03 1.6228E-03 2.7310E-03 5.3118E-02 1.0867E-02 6.7659E-03 2.7613E-02 8.8120E-03 1.2303E-02 1.9640E-02 317E-02 1.7327E-01 2.3892E-01 3.1576E-01 UP 27

473E-03 1.8626E-03 3.1360E-03 6.1262E-02 1.2638E-02 7.7938E-03 3.2297E-02 1.0353E-02 1.4466E-02 2.3115E-02 148E-02 2.0483E-01 2.8348E-01 3.7513E-01

567E-02 2.9377E-03 4.9381E-03 9.7120E-02 2.0095E-02 1.2281E-02 5.1312E-02 1.6495E-02 2.3054E-02 3.6870E-02 3.79E-02 3.2796E-01 4.5604E-01 6.0473E-01

143E-02 5.9136E-03 9.9323E-03 1.9845E-01 4.1512E-02 2.4807E-02 1.0682E-01 3.4688E-02 4.8571E-02 7.7840E-02 946E-02 6.9941E-01 9.8528E-01 1.3142E 00 UP 30

229E-02 1.6848E-02 2.8312E-02 5.8143E-01 1.2355E-01 7.1287E-02 3.2352E-01 1.0687E-01 1.5022E-01 2.4141E-01

716E-01 2.1999E 00 3.1600E 00 4.2528E 00

UP 31

680E-02 8.2320E-03 1.3867E-02 2.8913E-01 6.2565E-02 3.5202E-02 1.6671E-01 5.5747E-02 7.8595E-02 1.2655E-01 728E-02 1.1648E 00 1.6910E 00 2.2854E 00 UP 32

377E-02 2.0025E-02 3.3725E-02 7.1494E-01 1.5758E-01 8.6243E-02 4.2420E-01 1.4327E-01 2.0237E-01 3.2648E-01 27E-01 3.0340E 00 4.4456E 00 6.0297E 00 UP 33

111E-01 2.1745E-02 3.6700E-02 7.9403E-01 1.8108E-01 9.5099E-02 4.9720E-01 1.7040E-01 2.4133E-01 3.9037E-01 620E-01 3.6744E 00 5.4416E 00 7.4081E 00 UP 34

015E-02 1.1466E-02 1.9438E-02 4.2918E-01 1.0129E-01 5.1076E-02 2.8454E-01 9.8999E-02 1.4071E-01 2.2814E-01 062E-01 2.1749E 00 3.2555E 00 4.4481E 00 UP 35

773E-02 5.1077E-03 8.6936E-03 1.9476E-01 4.7085E-02 2.3071F-02 1.3447E-01 4.7287E-02 6.7397E-02 1.0944E-01 894E-02 1.0526E 00 1.5871E 00 2.1739E 00 1 UP 36

4 6 .

396E-02 9.3219E-03 1.5976E-02 3.6519E-01 9.1829E-02 4.3045E-02 2.6870E-01 9.5926E-02 1.3724E-01 2.2332E-01 823E-01 2.1742E 00 3.3088E 00 4.5451E 00 UP 37

645E-02 645E-02 3.0668E-03 5.3262E-03 1.2473E-01 3.3543E-02 1.4691E-02 1.0184E-01 3.7120E-02 5.3380E-02 8.7108E-02 162E-02 8.6163E-01 1.3254E 00 1.8261E 00 UP 38

282E-03 1.8181E-03 3.2218E-03 7.6355E-02 2.2341E-02 9.1098E-03 7.0614E-02 2.6252E-02 3.7934E-02 6.2070E-02 919E-02 6.2309E-01 9.6647E-01 1.3342E 00 UP 39

626E-03 6.2040E-04 1.1468E-03 2.6176E-02 8.8095E-03 3.3236E-03 2.9359E-02 1.1138E-02 1.6168E-02 2.6548E-02 081E-02 2.7180E-01 4.2449E-01 5.8657E-01 TICALITY

MULTIPLICATION FACTOR NANCE ESCAPE PROBABILITY 1.041061E 00 8.337727E-01 RMALUTILISATIONFACTOR9.477757E-01RMALFISSIONFACTOR1.307422E00RMALMULTIPLICATIONFACTOR1.239142E00THERMALMULTIPLICATIONFACTOR9.284109E-02INITEMULTIPLICATIONFACTOR1.091651E00TA-28,INITIALCONVERSIONRATIOANDRELATIVECONVERSIONRATIO

.

	DELTA-28	ICR	RCR	
123456789011234 111134 	1.1301E-01 0.0 7.7379E-02 0.0 0.0 0.0 0.0 0.0 0.0 0.0 0.0 0.0 0	1.1853E 00 0.0 1.0551E 00 0.0 0.0 0.0 0.0 0.0 0.0 0.0 0.0 0.0	6.1022E-05 0.0 5.5043E-04 0.0 0.0 0.0 0.0 0.0 0.0 0.0 0.0 0.0 0	
L	0.129/E-02	1.0694E 00	4.00062-04	

'RON BALANCE FOR INFINITE REACTOR

,	LEAKAGE	ABSOR .	PROD.	SCATT.OUT	FLUX	FLUX*VOL.
,	LEAKAGE 6.7889E-03 7.8010E-03 1.1002E-02 5.9507E-03 1.30027E-03 1.3825E-03 2.7265E-03 1.39056E-03 1.39056E-03 1.39056E-04 1.3521E-04 1.3521E-04 1.3521E-04 1.3521E-05 5.15883E-04 1.3585E-05 5.25897EE-04 1.3585E-05 5.25897EE-04 1.3585E-05 5.25897EE-04 1.58568EE-04 1.58568EE-04 1.58568EE-04 3.57568EE-04 1.48448E-03 3.66669EE-03 3.66669EE-03 3.66669EE-03 3.66669EE-03 3.66669EE-03 3.66669EE-03 3.66669EE-03 3.66669EE-03 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.66669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669EE-02 3.6669	ABSOR. 3.8604E-02 2.0470E-02 1.1702E-02 2.4328E-02 1.68662E-02 2.4931E-02 3.99564E-02 3.99564E-02 3.054866E-04 7.4052E-04 3.1554836E-04 7.4052E-04 3.4574E-04 3.4574E-04 3.4574E-04 3.4574E-04 4.6800E-03 5.116802E-04 4.68002E-04 3.4575E-04 4.68002E-03 5.116802E-03 5.116802E-03 5.116802E-03 5.116802E-03 5.118802E-03 5.118802E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00 1.63042E-00	PROD. 1.0176E-01 4.6419E-03 0.0 0.0 0.0 2.1528EE-03 7.0021E-03 1.62325E-04 8.3532E-04 8.3532E-04 8.3532E-04 8.36532E-04 8.36532E-04 8.36532E-04 4.1051E-04 3.99907E-04 4.9896E-04 4.9896E-04 4.9896E-04 4.9896E-04 4.9896E-04 4.9896E-04 4.9896E-04 4.9897E-03 5.9097E-03 5.9097E-03 1.1665E-022 8.6036E-02 1.3507E-01 1.3507E-01 1.3507E-01 1.3507E-02 1.5496E-02 2.9488E-02 1.666E-02 2.9488E-02 1.666E-02 2.9488E-02 2.9488E-02 2.9488E-02 2.9488E-02 2.9488E-02 2.9488E-02 2.94888E-02 2.9488E-02 2.94888E-02 2.94888E-02 2.94888E-02 2.94888E-02 2.94888E-02 2.94888E-02 2.94888E-02 2.94888E-02 2.94888E-02 2.94888E-02 2.948888E-02 2.948888E-02 2.948888E-02 2.948888E-02 2.948888E-02 2.948888E-02 2.948888E-02 2.948888E-02 2.948888E-02 2.948888E-02 2.948888E-02 2.948888E-02 2.948888E-02 2.9488888E-02 2.948888E-02 2.948888E-02 2.948888E-02 2.9488888-02 2.9488888-02 2.9488888-02 2.9488888-02 2.9488888-02 2.9488888-02 2.9488888-02 2.948888-02 2.9488888-02 2.9488888-02 2.9488888-02 2.9488888-02 2.9488888-02 2.9488888-02 2.9488888-02 2.9488888-02 2.9488888-02 2.9488888-02 2.9488888-02 2.9488888-02 2.9488888-02 2.9488888-02 2.9488888-02 2.9488888-02 2.9488888-02 2.94888888-02 2.94888888-02 2.94888888-02 2.94888888-02 2.9488888-02 2.948888888-02 2.94888888-02 2.94888888-02 2.94888888888888888888888888888888888888	SCATT.OUT 6.3289E-01 1.0174E 00 9.9130E-01 9.9130E-01 9.6171E-01 7.7941E-01 7.7941E-01 8.2105E-01 4.1358E-01 1.230E-01 1.230E-01 1.230E-01 1.1776E-02 3.30F-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-02 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00 3.4074E-00	FLUX 1.8341E02 2.1501E02 3.0762E-02 3.07664E-02 1.4964E-02 1.41830EE-02 5.6604E-03 1.41830EE-02 5.6604E-03 1.4194E-03 1.41797EE-04 5.7573EE-04 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.0250404 2.025040404 2.0	FLUX*VOL. 7.3364E 00 1.7024E 01 2.2535E 01 1.2306EE 01 2.2535EE 00 2.2535EE 00 2.2535EE 00 2.2642E 00 8.1025E 00 2.2642E 00 8.4745EE-01 2.1392EE-01 1.1029EE-01 1.1239EE-01 1.14394EE-01 1.44993EE-01 1.44993EE-01 1.5377EE-01 1.1866EE 00 8.8861EE 00 1.1866EE 00 1.18685EE 00 1.1455EE
39	8.4652E-02	1.1002E 00	1.1906E 00	3.5358E 01	2.7294E-01 4.9898E-01	1.0918E 02 1.9959E 02

-

48 1

1

-GROUPS PARAMETERS FOR INFINITE REACTOR

-AVERAGED AD GROUP FINE GROUPS DIFF.COEFF. ABSORPTION PRODUCTION SCATT.OUT 1- 9 10-39 1.3460E 00 1.0197E 00 9.2736E-03 5.3564E-02 4.8012E-03 6.6893E-02 2.4874E-03 1 0.0 **NSFER KERNEL BETWEEN BROAD GROUPS** 2 FROM 1 Э 1 2.7565E-01 2 2.4874E-03 3.0522E-01 ERATOR-AVERAGED AD GROUP FINE GROUPS DIFF.COEFF. ABSORPTION PRODUCTION SCATT.OUT 1.2972E 00 8.7066E-01 3.5712E-05 7.2370E-05 1- 9 10-39 5.3444E-05 1.2285E-02 12 0.0 0.0 NSFER KERNEL BETWEEN BROAD GROUPS FROM 1 2 0 1 2.5150E-01 2 1.2285E-02 3.9848E-01 L-AVERAGED AD GROUP FINE GROUPS DIFF.COEFF. ABSORPTION PRODUCTION SCATT.OUT 1.3106E 00 8.9261E-01 2.5673E-03 7.9507E-03 1.3545E-03 9.8521E-03 9.6003E-03 1- 9 12 10-39 0.0 NSFER KERNEL BETWEEN BROAD GROUPS FROM 1 2 Π

1 2.5812E-01 2 9.6003E-03 3.8474E-01 5

.

'RON BALANCE FOR FINITE REACTOR

2.2954E 000 1.3052E 01 3.4407E 01 2.1398E 02 6.2012E 00 2.4805E 03 3.6477E 00 7.3555E 00 1.6678E 00 3.6554E 02 1.5291E 01 6.1166E 03 3.6477E 02 1.5297E 01 7.9899E 03 3.6477E 02 1.9975E 01 7.9899E 03 3.6477E 00 8.5739E 00 0.0 3.5147E 02 1.9975E 01 7.9899E 03 3.6477E 00 3.5147E 02 2.10943E 00 2.4101E 03 4.8463E 00 5.9226E 00 0.0 3.5847E 02 5.2541E 00 2.4101E 03 4.8463E 00 5.9226E 00 0.0 3.5847E 02 2.55192E 00 1.0077E 03 4.8463E 00 2.4168E 00 2.55982E 02 2.3039E 00 2.821E 03 4.2343E 01 1.2420E 00 6.6068E 01 1.8837E 02 2.3039E 00 2.8251E 03 4.2343E 01 1.2420E 00 6.6068E 01 1.6683E 02 2.3039E 00 9.2156E 02 4.9608E 00 1.3797E 01 3.6008E 01 1.7702E 01 3.2708E 02 8.0164E 02 3.5968E 01 3.0400E 01 4.6646E 01 2.2854E 01 1.7702E 01 3.6039E 02 2.6486E 01 3.0400E 01 4.6646E 01 1.9637E 01 3.2708E 02 8.0164E 02 3.5968E 01 3.0400E 01 4.6646E 01 1.9637E 01 4.0879E 01 3.6039E 02 2.6486E 01 1.4848E 01 1.22502E 01 1.0071E 01 7.8525E 01 3.8779E 02 1.2268E 01 1.4848E 01 1.23502E 01 1.0077E 01 5.5198E 02 3.6039E 02 1.2268E 01 1.4848E 01 1.23502E 01 1.0077E 01 5.5198E 01 3.8779E 02 1.2268E 01 1.4848E 01 1.35746E 01 5.5198E 01 3.8679E 02 1.5285E 01 1.6787E 01 1.6648E 01 1.35746E 01 6.7481E 01 2.5988E 02 3.3206E 01 1.48479E 01 1.66881E 01 1.35746E 01 6.7497E 01 3.66667E 02 1.66864E 01 1.38596E 01 3.8879E 02 3.3206E 01 1.4885E 00 1.16678E 02 1.88646E 00 4.41850E 00 3.6667E 02 1.4695E 00 1.6787E 01 1.66881E 00 1.135746E 00 4.0286E 02 3.3820E 01 1.4655E 00 2.1688E 00 1.168860E 01 1.33796E 00 5.7998E 02 3.3820E 01 1.4655E 00 2.26455E 00 1.16678E 02 1.8744E 00 7.4976E 02 3.6667E 02 1.8744E 00 7.4976E 02 3.6676E 02 1.8744E 00 7.4976E 02 3.6061E 02 3.32691E 00 1.4885E 00 1.6687E 02 1.8744E 00 7.4976E 02 3.6061E 02 3.32691E 00 3.9878E 00 1.10072E 00 3.8679E 00 3.6075E 00 3.8679E 00 3.7875E 00 3.8679E 00 3.8679E 00 3.8679E 00 3.8679E 00 3.6061E 02 3.7977E 00 3.8679E 00 3.7875E 00 3.8679E 00 3.8	LEAKAGE	ABSOR.	PROD.	SCATT.OUT	FLUX	FLUX*VOL.
1.28955 01 2.8951E 02 3.5843E 02 9.3803E 03 9.1796E 01 3.6718E 04	2.2954E 00 2.8028E 00 1.4340E 00 2.9072E 00 1.901563E 00 1.84393E 00 1.84393E 00 1.33633E 00 1.33633E 00 1.33633E 00 1.33633E 00 1.33633E 00 1.33633E 00 1.33633E 00 1.336339E 00 2.50564E 00 1.30504E 00 1.80399E 00 1.80395E 00 1.80395	1.3052E 01 7.3545E 00 4.1922E 00 1.35739E 00 5.9226E 00 6.1739E 00 5.9226E 00 6.1739E 00 1.3917E 01 1.2420E - 01 1.2420E - 01 1.25998E - 01 1.2508E - 01 1.2508E - 01 1.252845E - 01 1.26477E	3.4407E 01 1.6678E 00 0.0 0.0 0.0 7.5468E-01 2.4388E 00 6.6068E-01 3.0103E-01 3.1024E-01 3.1024E-01 1.5213F-01 1.4848E-01 1.4848E-01 1.4848E-01 1.4848E-01 1.4848E-01 1.4848E-01 1.4848E-01 1.4848E-01 1.4848E-01 1.4848E-01 1.4848E-01 1.4848E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4858E-01 1.4	2.1398E 02 3.6554E 02 3.3470E 02 3.5147E 02 3.5147E 02 2.7370E 02 2.8597E 02 1.8284E 02 2.8597E 02 1.6833E 02 2.8597E 01 1.66545E 01 1.19450E 01 1.23645E 01 1.23645E 01 1.23645E 01 1.66384E 01 1.66384BE 01 1.66384BE 01 1.66384BE 01 1.66384BE 01 1.66385E 02 1.66385E 01 1.66385E 01 1.66385E 01 1.66385E 01 1.66385E 01 1.66385E 02 1.66385E 02 1	6.2012E 00 1.5291E 01 7.7028E 00 1.9843E 01 5.2541E 00 2.5192E 00 2.5192E 00 2.3039E 00 3.17705E-01 2.2854E-01 1.9631E-01 1.0071E-01 1.0220E-01 1.0376E-01 1.3574E-01 1.3574E-01 1.3574E-01 1.3574E-01 1.3574E-01 1.3574E-01 1.3574E-01 1.4891E 00 1.8744EE 00 1.8744EE 00 1.8744EE 00 1.8744EE 00 1.8744EE 00 1.66637EE 00 3.0752EE 00 9.6013EE 00 3.0752EE 00 9.6013EE 00 1.5774E 00 2.7134E 00 1.1791E 00 2.1493E 01 9.1796E 01 9.1796E 01	2.4805E 03 3.0811E 03 3.0811E 03 3.08971E 03 2.4007763E 03 1.00763E 03 2.1016E 03 1.00763E 02 3.27402E 02 3.27402E 02 3.27402E 01 7.8225E 01 4.0879E 01 5.5198E 01 5.5198E 01 5.5198E 02 2.744185PE 01 5.55198E 02 2.744185PE 02 2.744185PE 02 3.2744185PE 02 3.2744185PE 02 5.5198E 02 3.2744185PE 02 3.2745185PE 02 3.274518500 3.274518500 3.274518500 3.274518500 3.274518500 3.274518500 3.274518500 3.274518500 3.274518500 3.2757518500 3.2757518500 3.2757518500 3.2757518500 3.2757518500 3.2757518500 3.2757518500 3.2757518500 3.2757518500 3.2757518500 3.2757518500 3.2757518500 3.2757518500 3.2757518500 3.2757518500 3.2757518500 3.2757518500 3.2757518500 3.2757518500 3.2757518500 3.2757518500 3.2757518500 3.2757518500 3.2757518500 3.2757518500 3.2757518500 3.2757518500 3.27575185000 3.27575185000 3.27575185000 3.27575185000 3.27575185000 3.27575185000 3.27575185000 3.27575185000 3.275751850000 3.27575185000000000000000000000000000000000

٠

2919

IRIAL BUCKLING - EIGENVALUE = 3.9349E-04
ICTIVE MULTIPLICATION FACTOR = 1.0003E 00
IRESSIVE TIME(SEC.) AT THE END OF MAIN = 382.490

- 50 -

.

2851	SYS1.CMPLIB	PASSED
2851	\$4569101.T125507.FP001.PINOCCHI.GOSET	PASSED
2851	SYSOUT	SYSDUT
2851	PUNCH1	PASSED
2851	VOL SER NOS- PUNCHI. LIRE	DELETED
2851	SYS69101.T125507.RP001.PINOCCHJ.LIBPIN	DELETED
285I 285I	VOL SER NOS= E0KSTI. SYS69101.T125507.RP001.PINOCCHI.AVSIND	DELETED
žgoi k	383, LIBPIN, PINUCCHI	
2851	SYS69101.1125507.RP001.PINOCCHI.GOSET	DELETED
2851	VOL SER NOS- EURSYS. SYS1.CMPLIB	KEPT
PINOC	CHI DATE=69.04.11 BEG.T.=17.018 DURATION=0.	.121 N.OPER=463

S P JOB STATISTICS -- 1,735 CARDS READ -- 965 LINES PRINTED --

١

1

0 CARDS PUNCHED -- 7.31 MINUTES EXECUTION TIME

t 51

1.

<u>Table</u> 7

- 52 -

Isotope in HEROIC Library

ELEM.	1st IDENT.	2nd IDENT.	MODEL
Н	10001	27	FREE GAS
	10001	40	30 10
	10001	5 0	81 8 1
	10001	60	10 11
	10001	70	10 BT
	10001	80	11 11
	10001	9 0	¥¥ ¥¥
	12001	27	H in H ₂ 0 without transp.correction
	14001	27	H in H.O with
		-,	transp.correction
	12001	70	H in H O
	12001	150	" "2"
	12001	250	11 IT
	12001	350	88 8 8
	13001	27	ARDENTE MODEL
	13001	100	20 80
	13001	150	Vê 19
	13001	200	89 88
	13001	25 Ú	19 9
	13001	300	8 8 88
	13001	350	88 88
	13001	400	89 98
	13001	450	17 FI
D	10002	27	FREE GAS
	10002	50	88 88
	10002	90	tt 19
	10002	100	20 P2
	10002	150	8.0 8.0
	10002	200	88 88
	10002	250	87 8 8
	10002	300	19 1 9

ELEM.	1st IDENT.	2nd IDENT.	MO	DEL
D	12002	27	NELKI	N MODEL
	12002	40	(from GAKER	101 groups)
	12002	50	11	11
	12002	6 0	**	**
	12002	70	11	Ħ
	12002	80	**	11
	12002	90	11	88
	14002	27	(from GAKER	30 groups)
	14002	40	•1	11
	14002	50	**	16
	14002	60	11	11
	14002	70	**	11
	14002	80	**	11
	1 400 2	90	**	11
С	10012	27	FREE	GAS
	10012	100	**	11
	10012	150	**	11
	10012	177	**	11
	10012	200	81	11
	10012	250	**	99
	10012	300	11	**
	10012	350	**	t #
	10012	400	11	11
	10012	450	11	91
	10012	620	98	11
	10012	627	••	11
	10012	700	**	**
	10012	720	99	88
	10012	800	**	17
	10012	820	ч,	**
	10012	900	17	**
	10012	928	T†	**
	10012	1000	11	**
	10012	1100	11	**
	10012	1200	**	17
	10012	1300	••	F1
	10012	1400	99	**
	10012	1500	89	71
	15012	27	PARKS	KERNEL

ELEM.	Ist IDENT.	2nd IDENT.	MODEL
0	10016	27	FREE GAS
	10016	-10	19 19
	10016	50	£1 ED
	10016	60	11 11
	10016	70	41 11
	10016	80	11 11
	10016	90	17 88
	10016	150	44 82
	10016	177	¥\$ \$\$
	10016	25 0	11 11
	10016	3 50	¥1 11
	10016	627	11 11
	10016	7 00	¥\$ \$\$
	10016	300	¥? 17
	10016	900	¥7 ¥7
	10016	1000	ta et
	10016	1100	FF 18
	10016	1200	TF FT
	10016	1300	Ft 11
	10016	1400	19 13
	10016	1500	81 FT
Fe	26	27	
Al	27	27	
Zr	40	27	
Th232	232	27	
U235	235	27	
U236	236	27	
U238	238	27	
Pu239	239	27	
Pu240	240	27	
Pu240	102 40	177	Dopp.Broadenéd
	10240	327	88 88
	10240	627	85 B8
	10240	927	88 88
Pu241	241	27	
Xe	135	27	
Во	10	27	
Cr	52	27	
Min	55	27	

ELEM.	1st IDENT.	2nd IDENT.	MODEL
Ni	58	27	
Cu	63	27	
ND	93	27	
Cd	112	27	
Sn	119	27	
Pb	207	27	
Sm	149	27	
U233	233	27	
Pa233	533	27	
Np239	539	27	
1		•	
Fast and epithern	nal cross sections		
н	1	27	
ת ת	2	27	
Bo	10	27	
C	12	27	
0	16	27	
Μα	24	27	
Fe	26	27	
Al	27	27	
Zr	40	27	
Cr	52	27	
Mn	55	27	
Ni	58	27	
Cu	63	27	
Nb	93	27	
Cd	112	27	
Sn	119	27	
Хе	135	27	
Sm	149	27	
Pb	207	27	
Th	232	27	
U233	233	27	
U235	235	27	
U236	236	27	
U238	238	27	
Pu239	239	27	

ELEM.	1st IDENT.	2nd IDENT.	MODEL
Pu240	240	27	
Pu241 Pa233	533	27 27	
Np239	539	27	
Resonance	integral coeff	icients	
Th	232	27	
U235	235	27	
U238	238	27	

HEROIC Calculation

12	16	24	32
7.3	8.9	12.3	17.0
12.1	15.4	21.9	33.0
21.9	28.4	41.5	65.0
42.0	52.0	75.0	107.1
	12 7.3 12.1 21.9 42.0	12167.38.912.115.421.928.442.052.0	1216247.38.912.312.115.421.921.928.441.542.052.075.0

Computing times in minutes

Fast and epithermal cross sections





Fig. III





References

- AMYOT L. et al.: "Heavy Water Lattice Analyses at Ispra", BNES Int. Conf. London June 1967 on the Physics Problems in Thermal Reactor Design - Paper 15 pg. 117-125
- 2. DE HAAN W. et al. :"PLUTHARCO, A Plutonium, Uranium, Thorium Assembly Reactivity Code ", EUR-3141.e. (1966)
- AMYOT L. et al.: "PINOCCHIO: A Computer Programme for Cell Reactivity Calculations", not available
- 4. CARLVIK I.: "A Method for Calculating Collision Probabilities in General Cylindrical Geometry and Applications to Flux Distribution and Dancoff Factors", Paper P/681, Third Geneva Conf. (1964)
- 5. VIEWEG M.A. et al.: "GGC-II Program for Using GAM-II and GATHER-II Spectrum Codes in Preparing Multigroup Cross-section input", GA-4436 (1963)
- 6. WESTCOTT C.H.: "A Survey of Values of the 2200 m/ Constants for Four Fissile Nuclides", Atomic Energy Review - Vol. 3 -No. 2 (1965)
- 7. ARDENTE V.: "Remarks on the Slow Neutron Scattering by Organic Molecules", J. de Physique, <u>25</u>, 64 (1964)
- 8. DE HAAN W.:"Correlation of the Analytical Theory of the Results of Nordheim Numerical Resonance Integral Calculation Method", Private Communication
- 9. BREVI R.: "Analisi Critica del Calcolo dell'Integrale di Risonanza del Torio" -
- 10. DE HAAN W.et al.: "Calculation of the Absorption and Fission Resonance Integrals of U-235 by a Fitting Procedure " EUR 3936e. (1968)

- 11. LESLIE D.C., JONSON A.: "The Calculation of Collision Probabilaties and Resonance Integrals in Cluster Type Fuel Elements", AEEW-k 384 (1964)
- 12. BEARDWOOD J.E. et al.: "The Jolution of the Transport Equation by Collision Probability Methods", Conference on the Application of Computing Methods to Reactor Problems, ANL-7050 pp. 93-110 (1965)
- 13. HONECK M.C.: "A Thermalisation Transport Theory Code for Thermal Lattice Calculations", BNL-5826 (1962)
- 14. BENOIST P.: "Streaming Effects and Collision Probabilities in Lattices", CEA Report SPM n. 988 (1967)
- 15. BOEUF A., TASSAN S.: "A Measurement of the Hyperfine Structure of the Thermal Neutron Flux Distribution in Organic Cooled UC Clustered Fuel Elements in a D₂O Moderated Lattice" - Energia Nucleare Vol. 13, n. 2 (1966)
- 16. KAY R.E., GREEN K.E.: "Lattice Measurements with 7-rod Clusters of Natural Uranium Carbide in Heavy Water Moderator", AECL-2650 Chalk River, Ontario (1966)
- 17. BOEUF A., TASSAN S.: "A Measurement of the Fine Structure Distribution of the Thermal Flux in Organic Cooled Fuel Elements in a D₂O Moderated Lattice", Energia N^CCleare Vol. 11, n. 10 (1964)
- 18. JONSSON A., PEKAREK H.: "Multigroup Collision Probability Theory in Cluster Geometry, Comparison with Experiments", Reactor Physics in the Resonance and Thermal Regions, San Diego pag. 367-393 (1966)
- 19. Misure di K-inf e di Parametri Dettagliati per Reticoli CIRENE ai Passi 27.4,30.2 e 32.2 cm e con Densità del Refrigerante Equivalente a 0.35 gr/cm³ di H_20 - Doc. Int. CISE n. 279

NOTICE TO THE READER

All Euratom reports are announced, as and when they are issued, in the monthly periodical "euro abstracts", edited by the Centre for Information and Documentation (CID). For subscription (1 year : US 16.40, £ 6.17) or free specimen copies please write to :

Handelsblatt GmbH " euro abstracts " Postfach 1102 D-4 Düsseldorf (Germany)

> Office de vente des publications officielles des Communautés européennes 37, rue Glesener Luxembourg

To disseminate knowledge is to disseminate prosperity — I mean general prosperity and not individual riches — and with prosperity disappears the greater part of the evil which is our heritage from darker times.

Alfred Nobel

SALES OFFICES

All reports published by the Commission of the European Communities are on sale at the offices listed below, at the prices given on the back of the front cover. When ordering, specify clearly the EUR number and the title of the report which are shown on the front cover.

SALES OFFICE FOR OFFICIAL PUBLICATIONS OF THE EUROPEAN COMMUNITIES

37, rue Glesener, Luxembourg (Compte chèque postal Nº 191-90)

BELGIQUE — BELGIË MONITEUR BELGE

MONITEUR BELGE rue de Louvain 40-42 - 1000 Bruxelles BELGISCH STAATSBLAD Leuvenseweg 40-42 - 1000 Brussel

DEUTSCHLAND

BUNDESANZEIGER Postfach - 5000 Köln 1

生活情

FRANCE

SERVICE DE VENTE EN FRANCE DES PUBLICATIONS DES COMMUNAUTES EUROPEENNES 26, rue Desaix - 75 Paris 15e

ITALIA

M IN IN

设行 印刷制件信用

LIBRERIA DELLO STATO Piazza G. Verdi, 10 - 00198 Roma

LUXEMBOURG

OFFICE DE VENTE DES PUBLICATIONS OFFICIELLES DES COMMUNAUTES EUROPEENNES 37, rue Glesener - Luxembourg ð

û

ð

NEDERLAND STAATSDRUKKERIJ Christoffel Plantijnstraat - Den Haag

UNITED KINGDOM H. M. STATIONERY OFFICE P.O. Box 569 - London S.E.1

> Commission of the European Communities D.G. XIII - C.I.D. 29, rue Aldringer Luxembourg